

Origins and probes of cosmic magnetic fields

– *an introduction*

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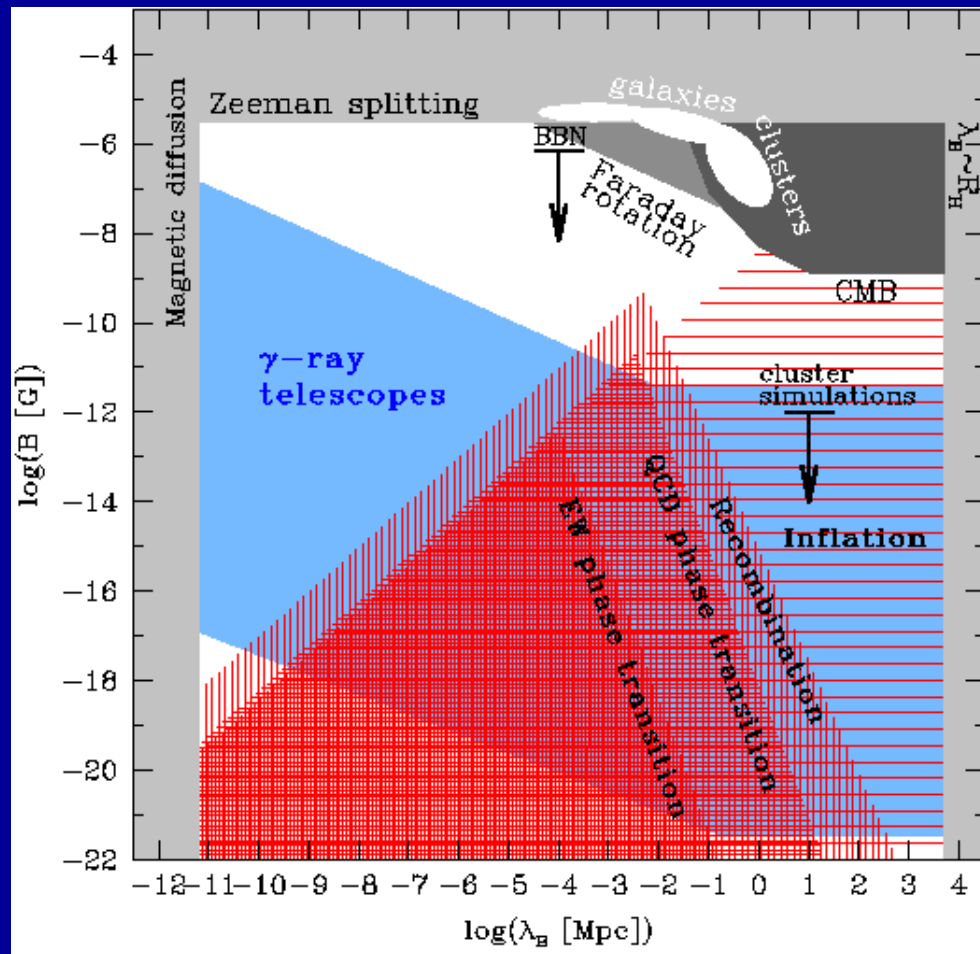
and

Los Alamos National Laboratory, USA

Paris, France, le 13 décembre, 2010

The CMF conference logo

D. Semikoz



Magnetic fields beyond the bounds of individual galaxies and clusters.

Overview

- This talk mostly excludes B within galaxies and within galaxy clusters
- Local universe, filaments, voids, recombination era. ($z \approx 0$ to $\gtrsim 1500$).
- opportunities for next-generation observations
- Some guiding theory: modelling: energy, physics
- Synergies *e.g.* with high energy- & astroparticle physics.

Comments on seed fields

Possibilities:

1. Primordial B
 - i.e. before the epoch of last scattering.
2. B seeded in baryonic plasmas, post-recombination, in stars/ galaxies/ IGM

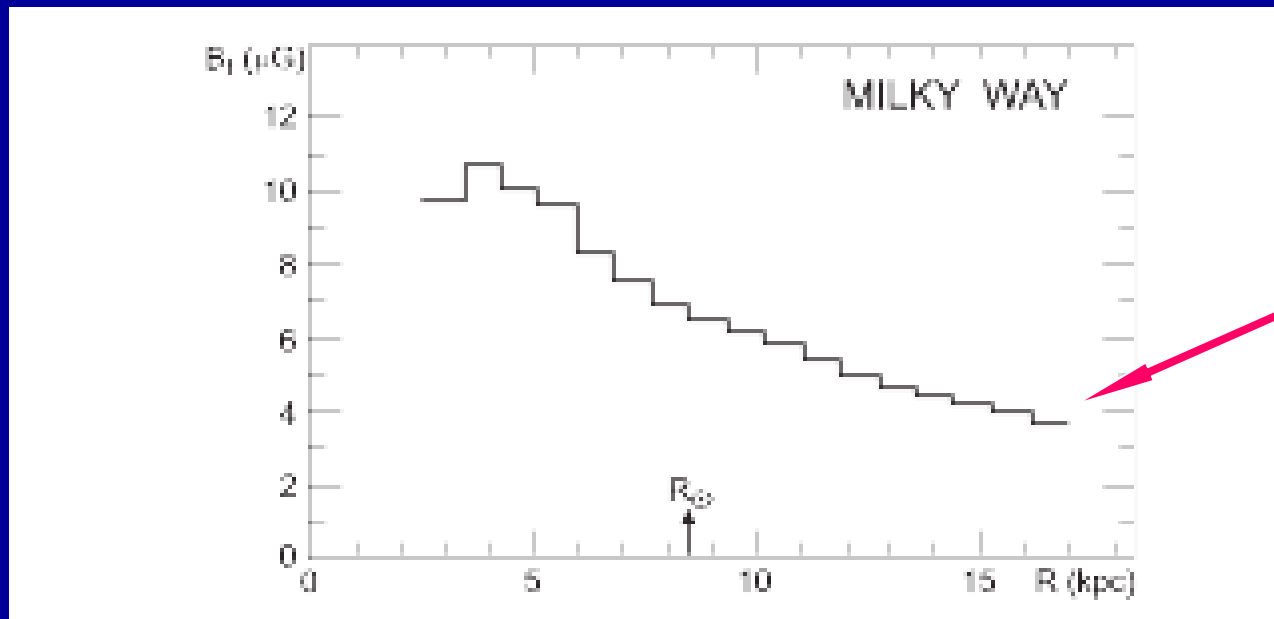
Comments:

- Possibility 1 difficult to verify – but
 - important connections to physics, particle cosmology, and possibly string theory.
 - Possibilities start at the Planck scale
 - Important test: $\langle |B| \rangle$ in voids
 - (*reviews: e.g. Kronberg 1994, Widrow 2002, Grasso & Rubenstein 2004*)
- Possibility 2 is certain: Mechanisms experimentally verified
 - star + SN-driven outflow
 - Supermassive BH – produced fields
 - Subsequent regeneration mechanisms.

1.

B in the vicinity of galaxies

Model of the Galactic $|B|$ vs. r .
from all-sky, 0.4 GHz synchrotron emissivity
(Haslam et al.)
supported by γ -ray observations
(Strong et al.)



Connects to
local B_{IGM} ?

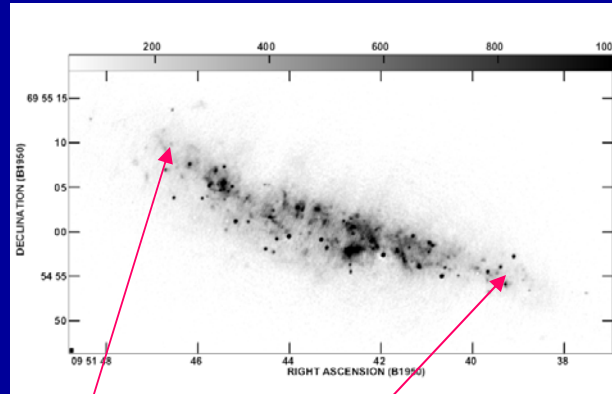
E.M. Berkhuijsen, R. Wielebinski (MPIfR Bonn)

How galactic winds
driven by stellar processes
inject magnetic fields into the IGM

- A. “quiescent”, Milky Way – like galaxies
- B. Starburst and dwarf galaxies

Outflow to IGM from the M82 starburst galaxy (3 Mpc distant)

Outflow halo:
 $B_{\text{halo}} \sim 10 \mu\text{G}$
coherence scale
1 – 2 kpc



VLA, All-config, 5 & 8 GHz, $\sim 0.3''$ resolution

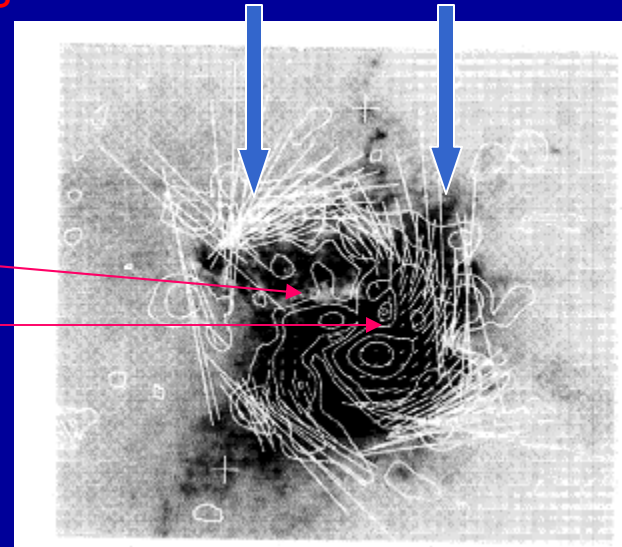
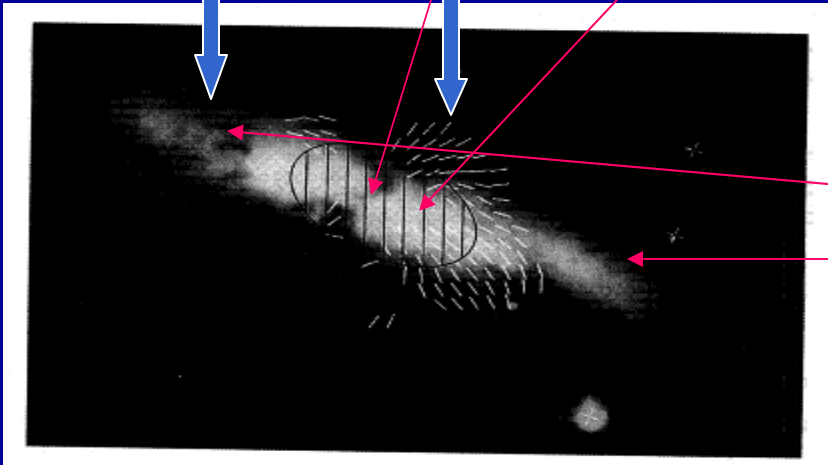
Kronberg, P.P. Biermann, P.L. Schwab, F.R. *ApJ* 246, 751, 1981.

Allen & Kronberg, 8GHz

Optical image

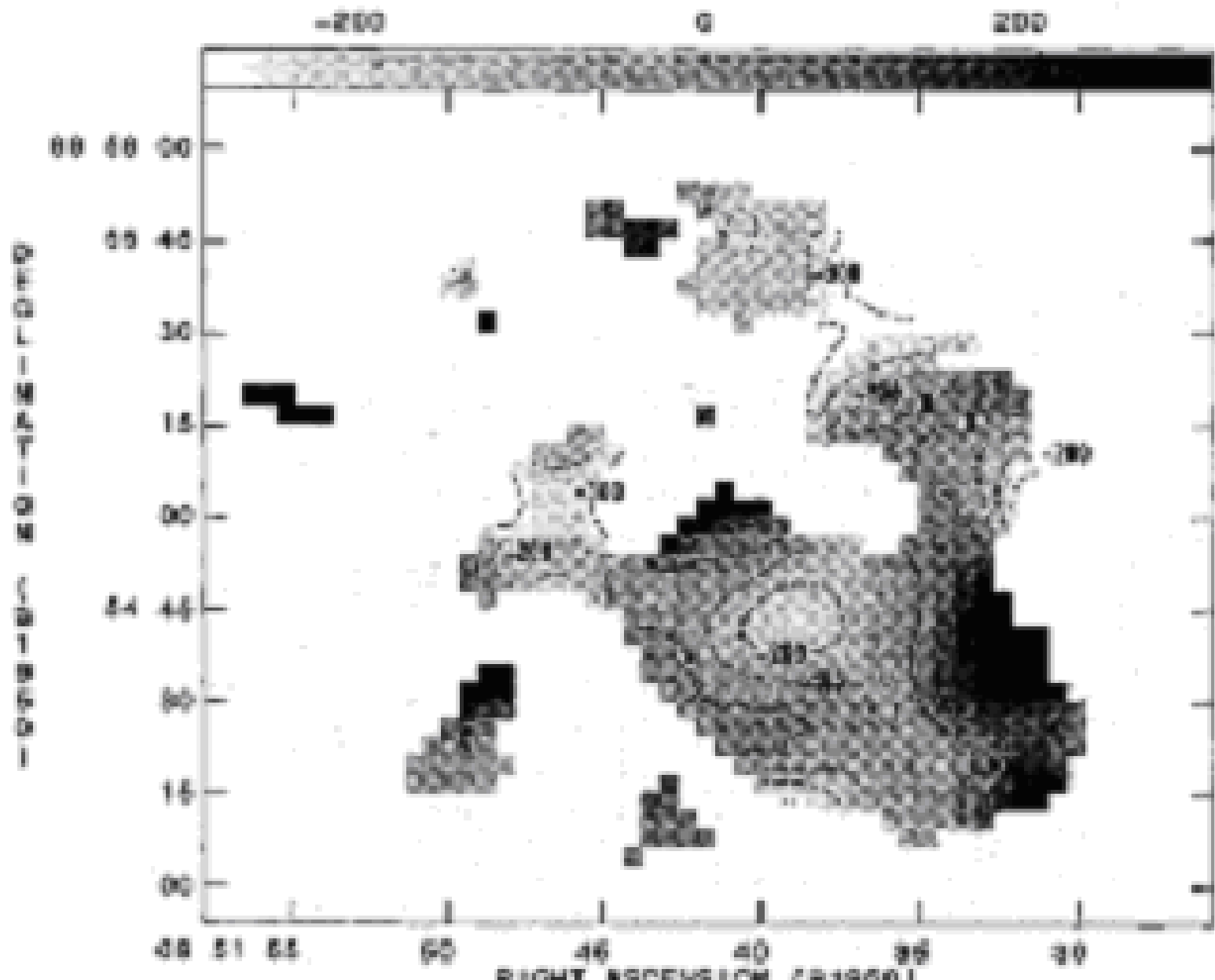
De-Faraday rotated,
projected magnetic field lines
From $\lambda\lambda$ 3.6 & 6.2 cm

Pol'n intensity $H\alpha$ emission



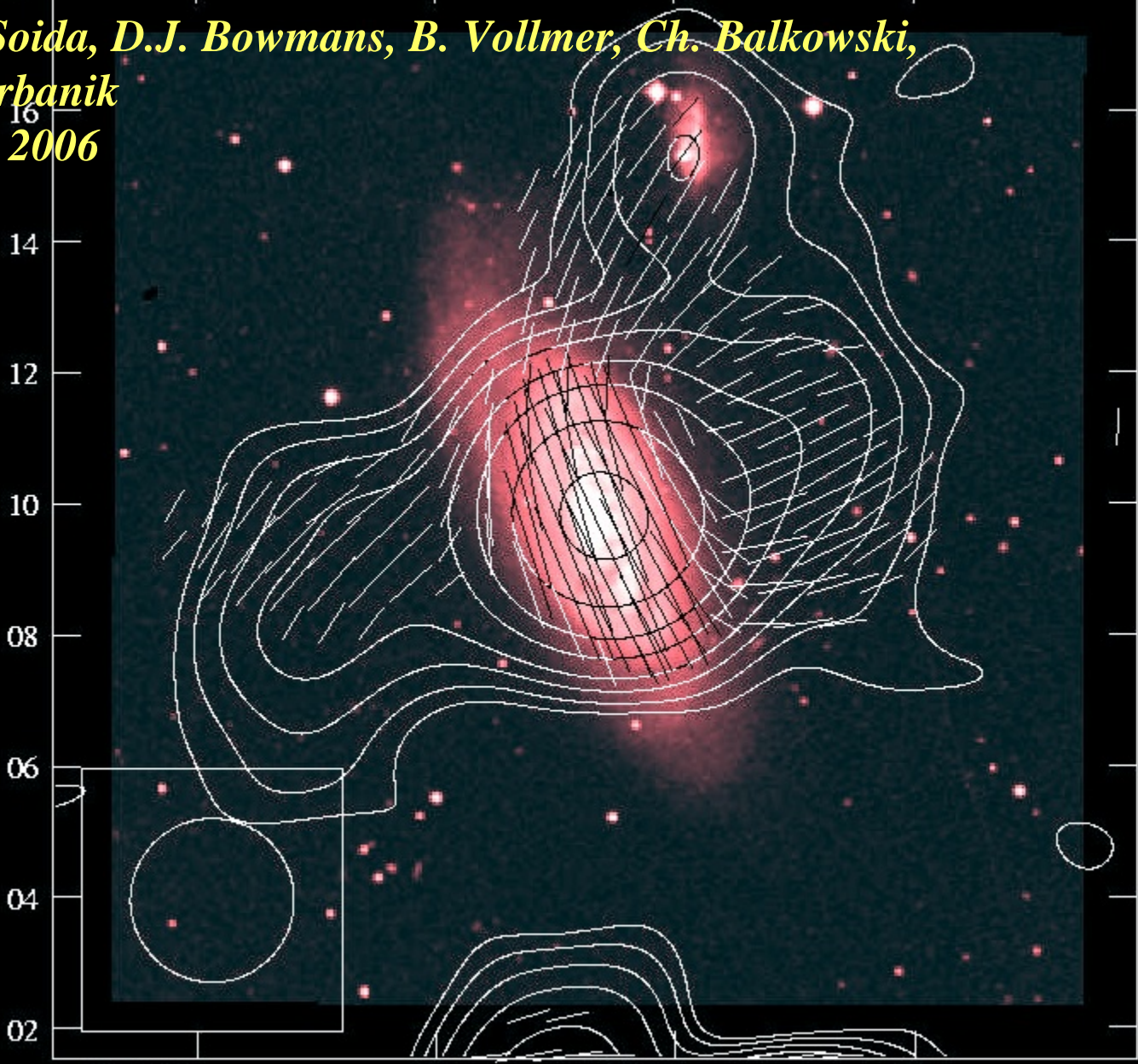
Reuter, H.-P., et al.. *A&A*, 282, 724, 1994, [*A&A* 293, 287, 1995 - Figs. with corrected orientation].

M82 Rotation Measure



13 18

DECLINATION (J2000)



02

04

06

08

10

12

14

16

12 37 15

00

36 45

30

RIGHT ASCENSION (J2000)

*K. Chyzy, M. Soida, D.J. Bowmans, B. Vollmer, Ch. Balkowski,
R. Beck, M. Urbanik
A&A 347,465, 2006*

2.

Collective dwarf galaxy seeding of the IGM MODEL INPUTS

- Outflow halo parameters at low z
- Dwarf galaxy counts, merging models
- Hierarchical merging scenarios since $z \sim 15$
- Embed in Hubble flow

RESULTS

- Volume of intergalactic filaments at $z \sim 0$
is easily filled with magnetic fields at $B_0 \sim 10^{-9}$
before any post-amplification --see below

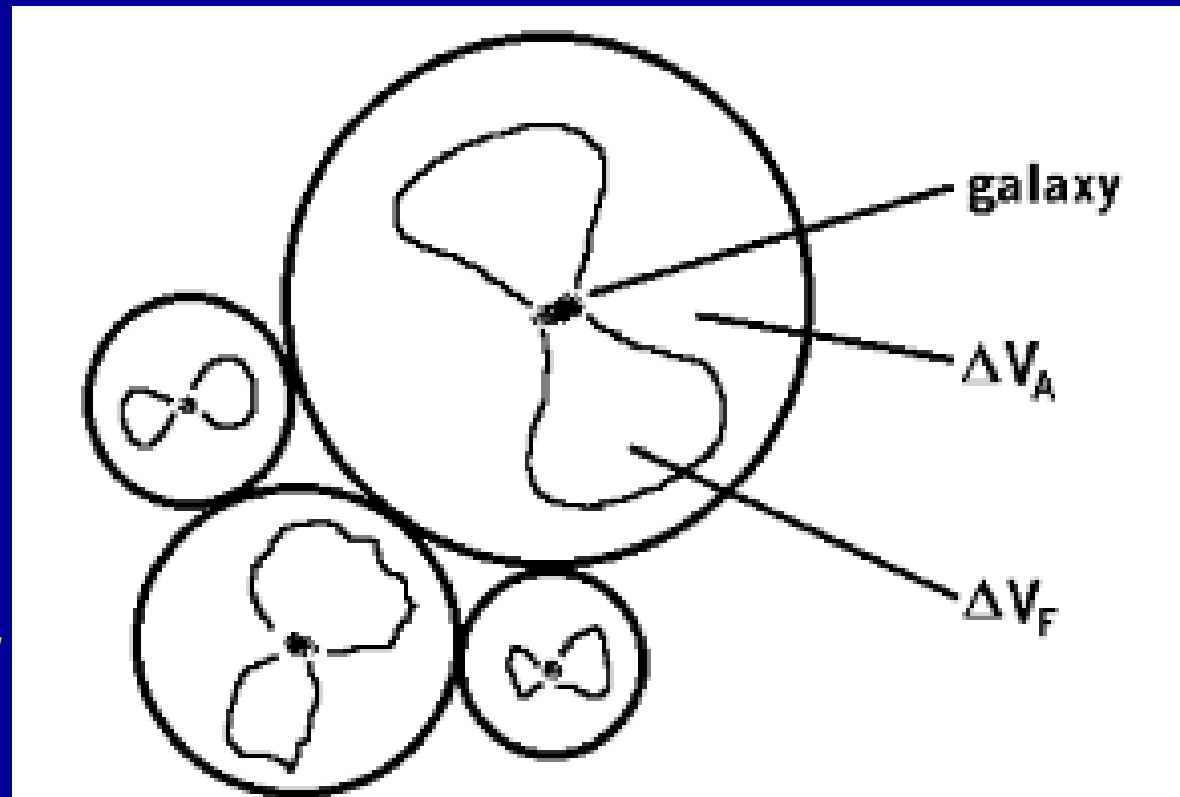
V_A = i.g. volume “**available**” to be filled with outflow

$V_F(t)$ = volume **filled** with stellar/SN halo outflow

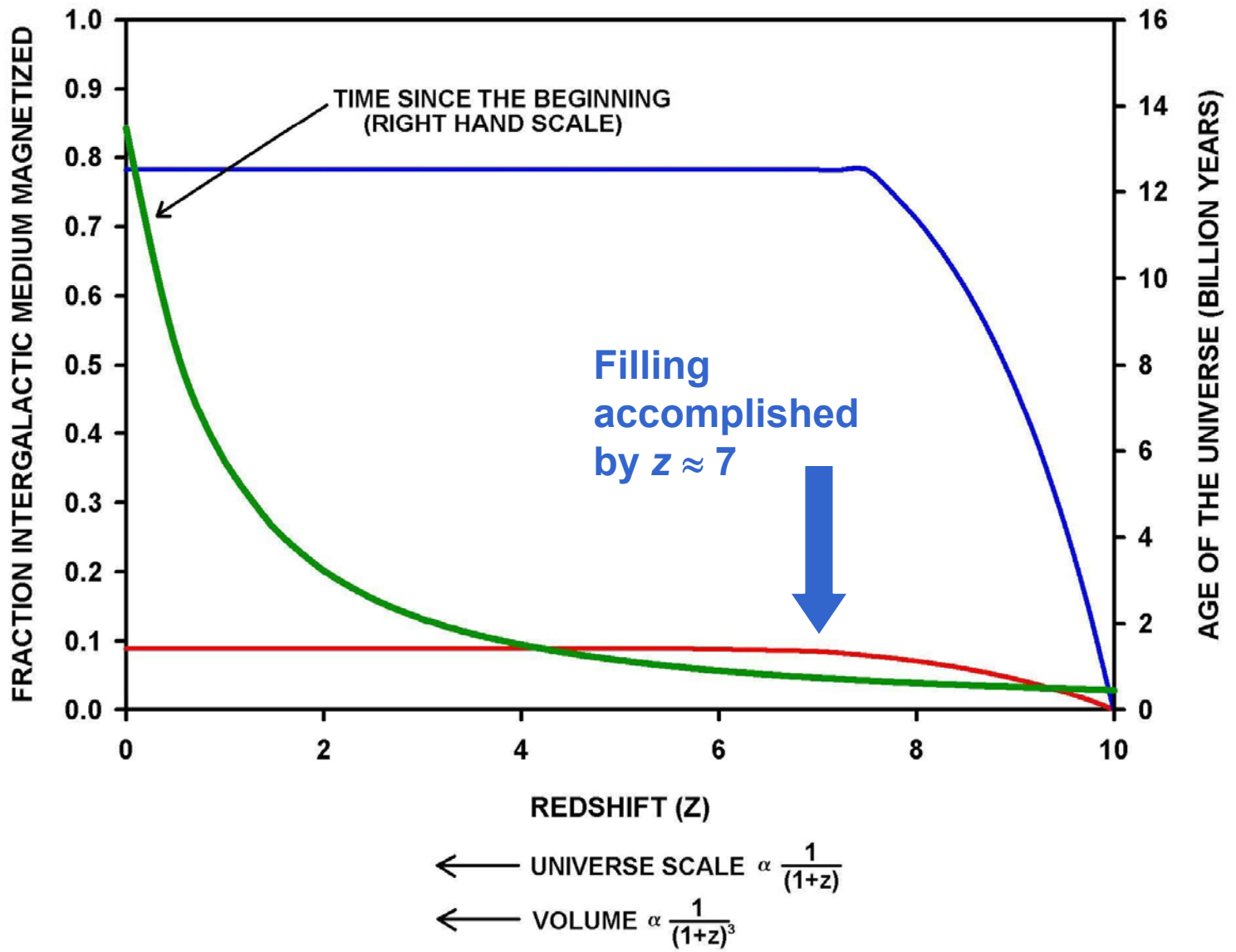
$f(t)$, [or $f(z)$] = V_F/V_A fraction of available IGM (*i.e.* within galaxy filaments) that gets filled from $z \sim 10$ to $z = 0$

An analytical model
for early starburst
dwarf galaxy seeding
of the IGM with
magnetic fields

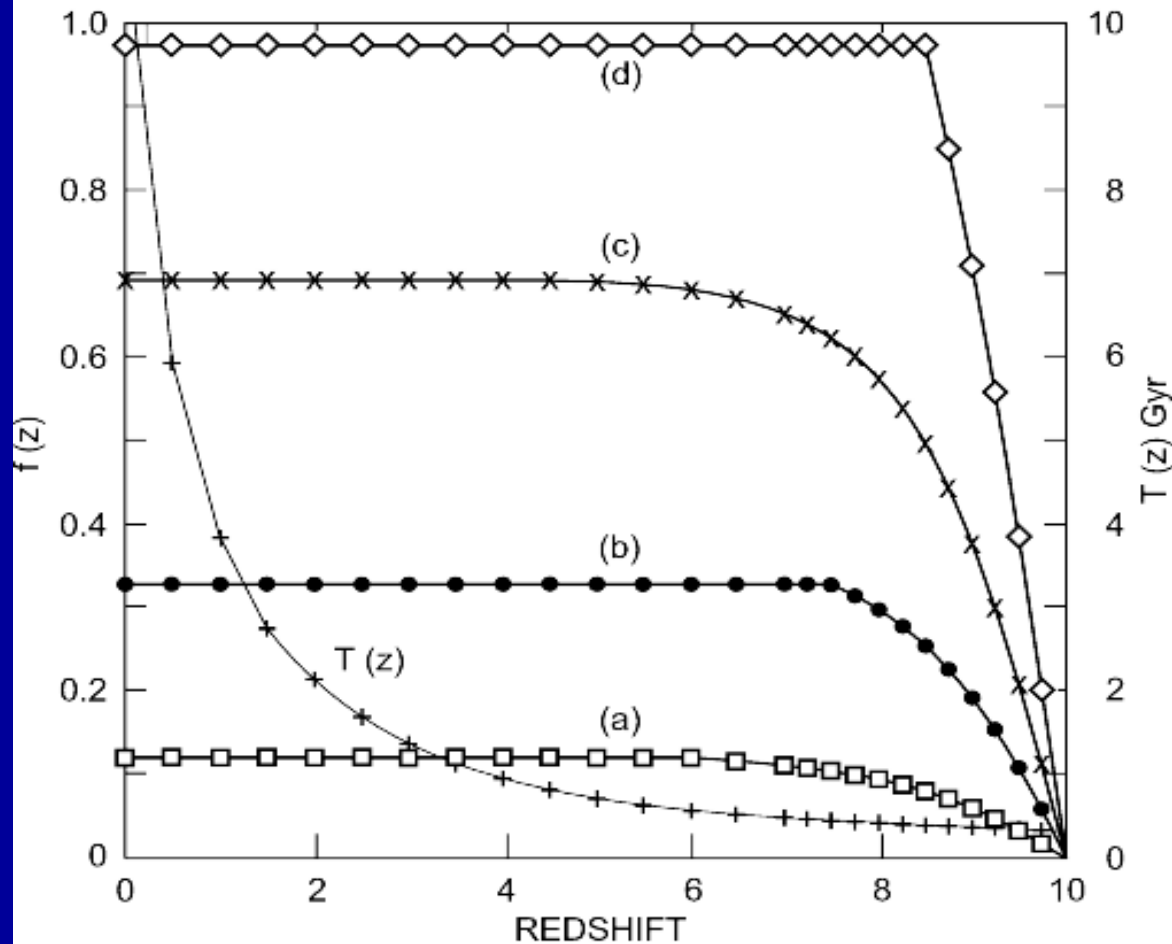
*P.P. Kronberg, H. Lesch,
and U. Hopp, ApJ, 511, 56-
64, 1999*



FILLING OF INTERGALACTIC SPACE BY EARLY STARBURST GALAXY OUTFLOWS



I.G.M. MAGNETIC FIELD VOLUME FILLING FACTOR



- 10Myr, $10 > z > 4$, $m=2$, $\rho_0 = 0.5$ $t=180$ Myr
- 10Myr, $10 > z > 7.5$ $m=3$, $\rho_0 = 0.5$ $t=100$ Myr
- x— 10Myr, $10 > z > 0$, $m=3.5$, $\rho_0 = 0.3$ $t=260$ Myr
- ◇— contin. $10 > z > 8.5$ $m=3.5$, $\rho_0 = 0.5$ $t= 73$ Myr
- +— Proper Time $T(z)$ Gyr

RESULTS & CONCLUSIONS

IGM MAG FIELD
SEEDED
by
STARBURSTING
PRIMEVAL GALAXIES

Figure from:
*P. Kronberg, H. Lesch & U.
Hopp ApJ, 511, 56-64, 1999*

A starting
template for
full simulations

Subsequent B - amplification of early galactic wind fields?

- $B_{i.g.}$ Subsequently amplified in vortices of shearing flows in LSS formation?
- up to $\sim 10^{-7}G$ in LSS filaments?

D. Ryu, H. Kang, & P. L. Biermann, A&A, 335, 19, 1998,

H. Kang, S. Das, D. Ryu, J. Cho, ICRC 2007,

D. Ryu, Kang, H, Cho, J., Das, S. Science, 320, 909 2008:

H. Kang, S. Das, D. Ryu, J. Cho
ICRC 2007 Proceedings, and
Science, 2008

earlier work: Ryu, Kang & Biermann ApJ 1998

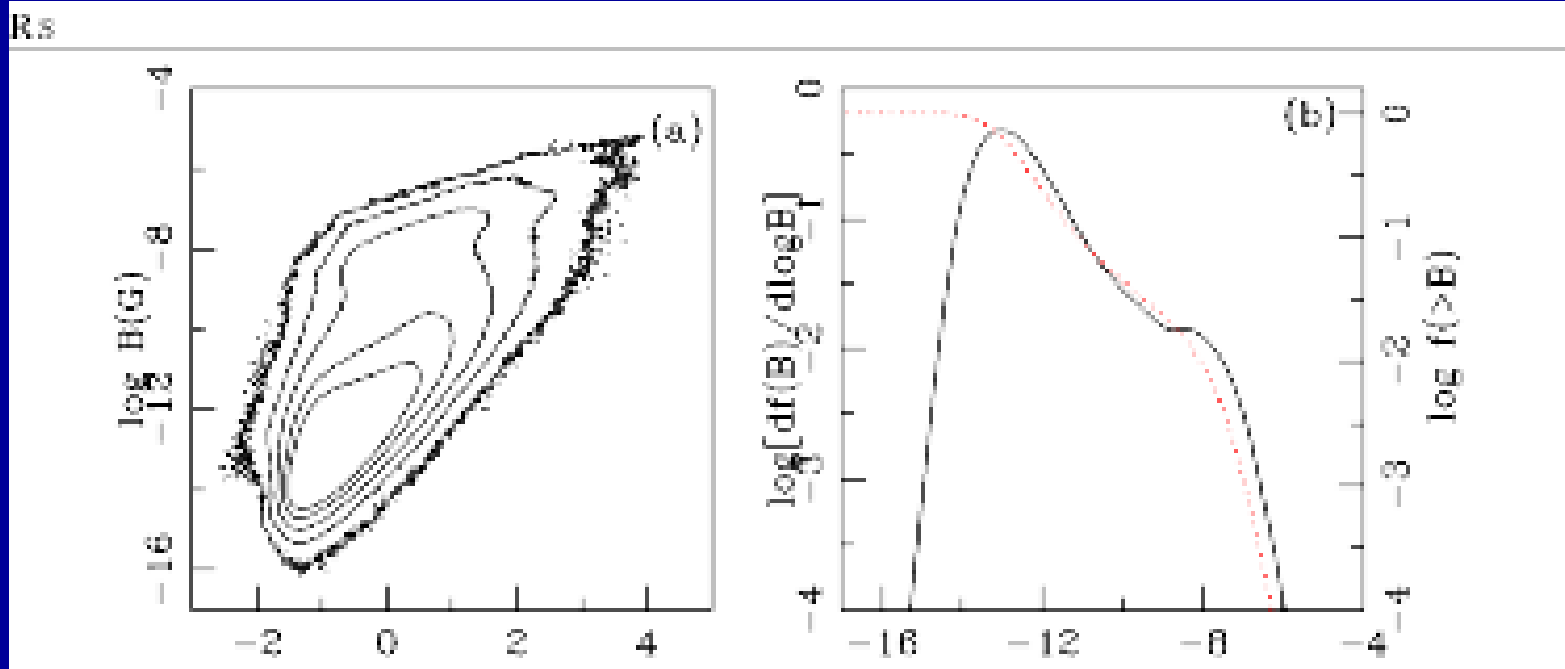


Figure 1: (a) Volume fraction in the gas density-EGMF strength plane with our model EGMF at $z = 0$. (b) Volume fraction, $df/d\log(B)$, (solid line) and its cumulative distribution, $f(>B)$, (dotted) as a function of EGMF strength.

3.

Central galactic black holes as a source of IGM magnetic energy

1. Energetics
2. Global consequences for IGM fields
3. Connections to fundamental plasma processes

IGM magnetic energy supplied by central galactic black holes

Can be globally quantified

A global, observation-based calculation:

Average
BH density
($M_{BH} \gtrsim 10^6 M_{\odot}$)

$$\langle \rho_{BH} \rangle \geq 2 \times 10^5 M_{\odot} / \text{Mpc}^3$$

Gravitational energy
reservoir per BH
(scaled for infall to R_G)

$$M_{BH} c^2 = 1.8 \times 10^{62} \frac{M_{BH}}{10^8 M_{\odot}} \text{ ergs}$$

This leads to an average magnetic energy density, ε_B
supplied to the IGM
from supermassive black holes .

if no B-dissipation over \sim a Hubble time

Smoothed out SMBH magnetic energy reservoir

$$\varepsilon_B = 1.36 \times 10^{-15} \left(\frac{\eta_B}{0.1} \right) \times \left(\frac{f_{RG}}{0.1} \right) \times \left(\frac{f_{FILAMENTS}^{VOL}}{0.1} \right)^{-1} \times \left(\frac{M_{BH}}{10^8 M_\odot} \right) \text{ erg cm}^{-3}$$

Gives $B_{IG}^{BH} = \sqrt{8\pi\varepsilon_B} = 1.8 \times 10^{-7} \text{ G}$

- Initially captured within galaxy filaments

Conclusion:

IGM near galaxies should contain magnetic energy $\approx \varepsilon_B$

• Next:

Observational tests for ε_B in the IGM

4.

Two recent probes for magnetic fields in local cosmological LSS, *beyond* galaxy clusters

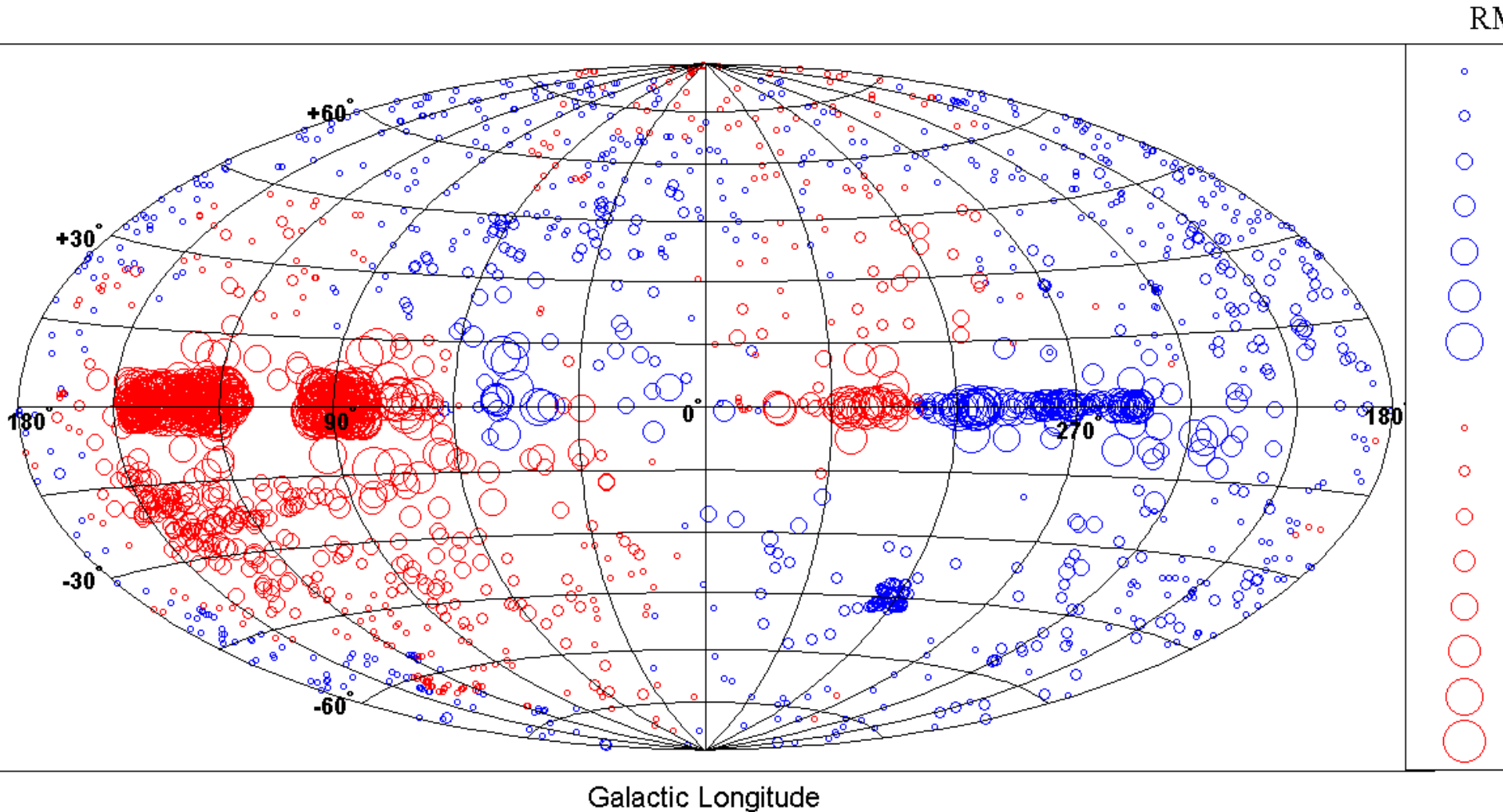
1. First Faraday RM probe for $\langle |B| \rangle$ in local universe LSS filaments of galaxies defined from (1) CfA2 and (2) 2MASS surveys

Xu, Kronberg, Habib & Dufton ApJ 637, 19, 2006

2. Search for unprecedented, faint, synchrotron radiation combining the 305m **Arecibo telescope** and the 1000m **DRAO interferometer**. Capable of $\sim 0.1 \mu\text{G}$ – level B_{IGM} detection.

Kronberg, Kothes, Salter, Perillat ApJ 659, 267, 2007

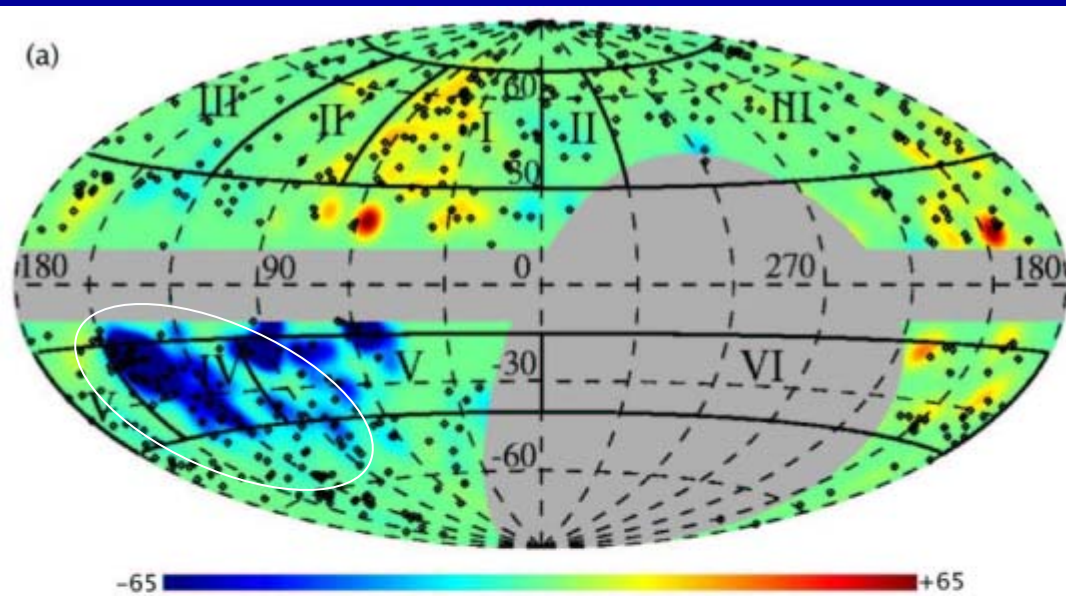
New smoothed Galactic RM sky from 2250 egrs RM's



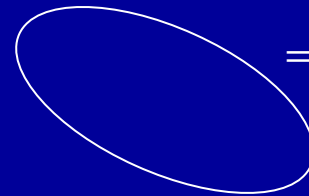
Probe No.1

RM search for magnetized plasma beyond clusters

Xu, Kronberg, Habib, Dufton: ApJ 2006, 637, 19

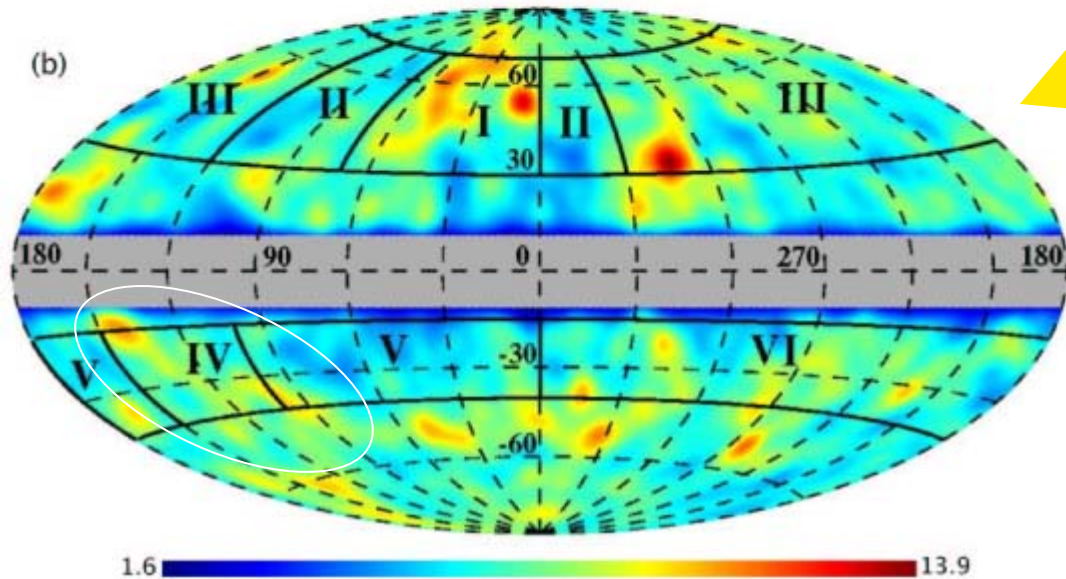


SMOOTHED
FARADAY
ROTATION



= Perseus-Pisces
supercluster

rad/m²



GALAXY COLUMN
DENSITY
(Method #2:
2MASS, HEALPix)



galaxies
per pixel
(\propto column density)



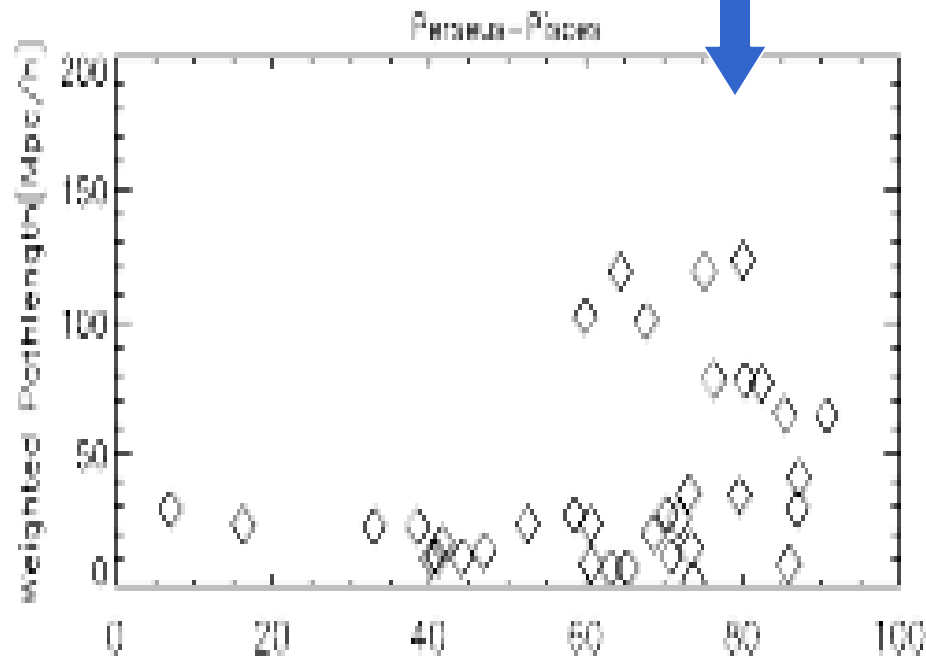
Optical galaxy counts vs. RM plots
for the Perseus-Pisces supercluster chain
Two types of investigation

Xu et al. ApJ 2006

(a)

Galaxy column density vs RM
from 7° -smoothed data

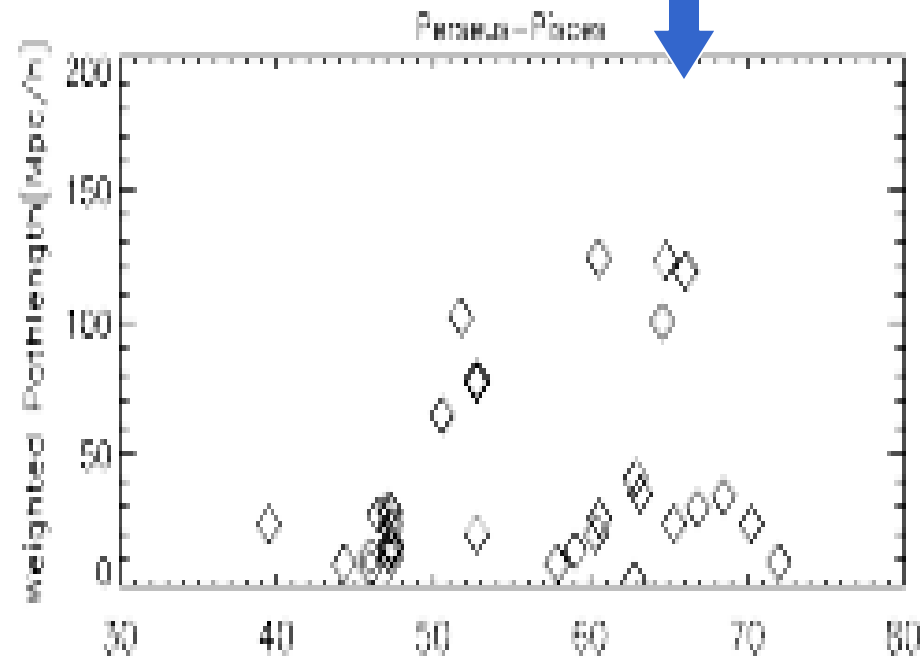
(used the 2MASS galaxy survey)



(b)

Weighted path length vs RM
from 3-D Voronoi-tessilated IGM filament volumes
(\because 3-D spectroscopic z 's are measured).
also from 7° -smoothed data

(Used the CfA2 galaxy survey)



Result (*Xu et al. ApJ 637, 19, 2006*)

- Attempted 3 local superclusters Virgo, Hercules, Perseus-Pisces. **2 independent galaxy survey analyses + RM's**
 1. CfA2 survey -spectroscopic z 's – 3D
 2. 2MASS survey spectral z 's – column densities
- Tentative result for B in Perseus-Pisces supercluster filament zones.
 $\sim 10^{-7}G$ using both CfA2 and 2MASS

Probe No. 2

B_{IGM} from diffuse
synchrotron radiation detection

A novel combination of :

The Arecibo radio telescope

(largest single radio reflector)

with

The DRAO interferometer

(Wide-angle, precision-imaging interferometer)

Astrophysical aims of Arecibo-DRAO radio images

- Detect weak magnetic fields in intergalactic space via diffuse synchrotron emission (at $< 1\text{GHz}$)
- Search possible radio **foregrounds to the cosmic microwave background** (CMB) on the scale of arcminutes (multipole scales, l , up to ~ 3000)
- Test for energy exchange from **central black holes of galaxies** to the **intergalactic medium**
- Can we detect a radio counterpart to the “Warm-hot intergalactic medium” (seen in soft X-rays)?
- Explore connections between the **radio**, and the **X- and gamma-ray** Universe)



Arecibo 305m Telescope, PR

*2 mm rms optics
illuminated area $\approx 200\text{m}$
uv overlap with DRAO $\approx 200\text{m}$*

Dominion Radio Astrophysical Observatory Penticton BC, Canada



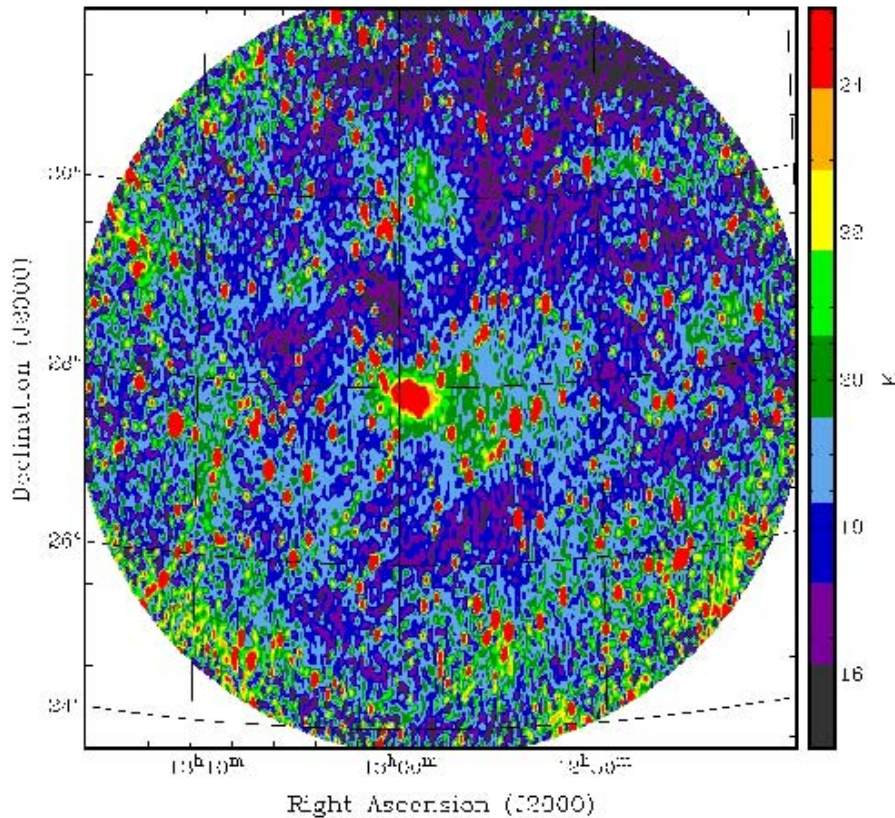
Precision track (E-W)

7 x 9m dishes

Max. separation = 617m \Rightarrow 1000m equiv. single dish resolution
Min. projected separation \approx 18m

In 12 days, 1 full image within 9° circle at 408 MHz

8° dia. Arecibo + DRAO
image, at a resolution of
2.5' x 6.5' 0.4 GHz

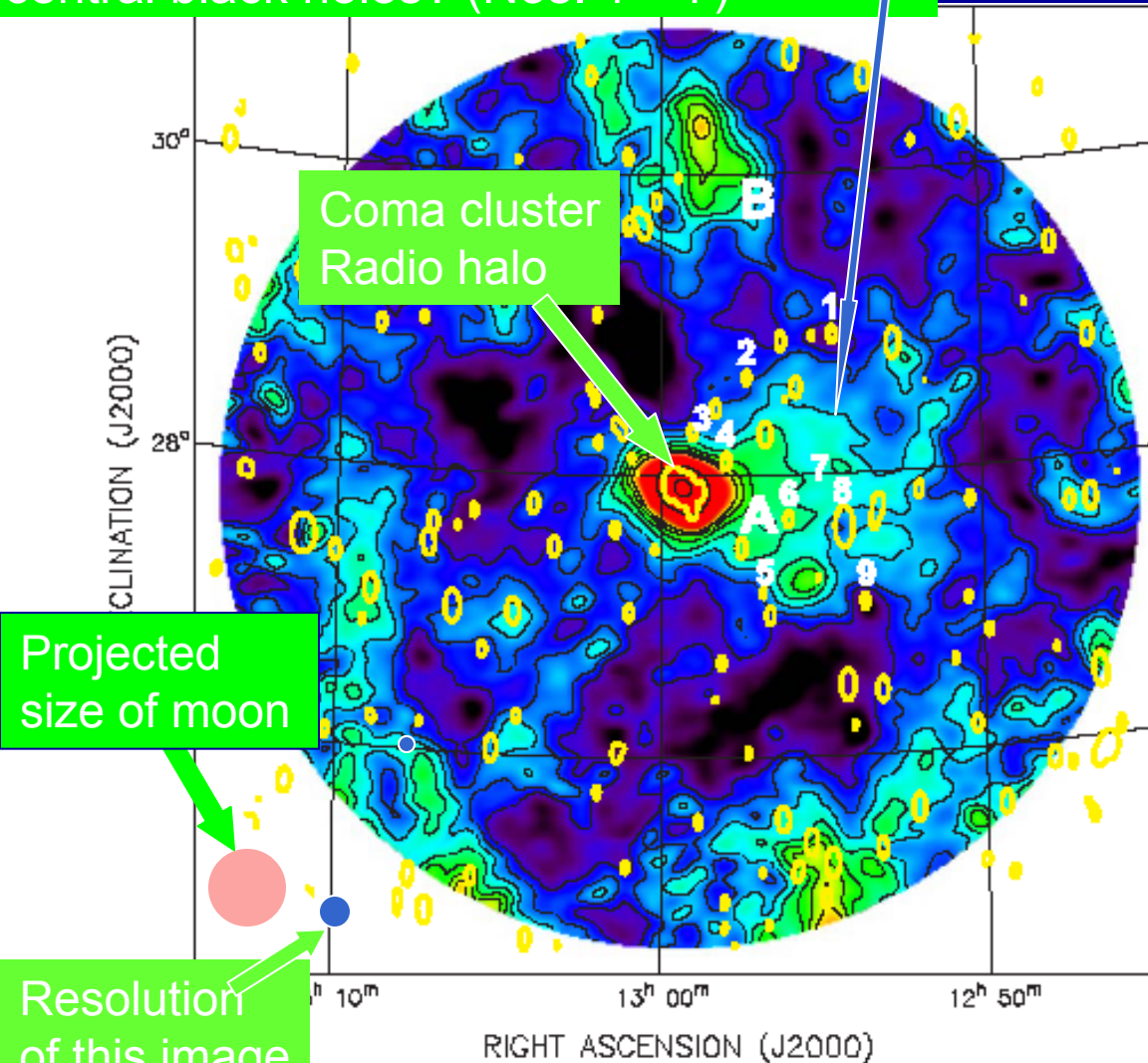


2.7K CMB background and
galactic foregrounds ($\approx 18K$)
are included

COMBINED Arecibo-DRAO image, now smoothed to **10'** (Arecibo) **resolution**

P. Kronberg, R. Kothes, C. Salter, & P. Perillat ApJ 659, 267, 2007

Collective energization of several galactic central black holes? (Nos. 1 – 7)



- Discrete sources removed,
- CMB + linear plane Milky Way foreground removed
- Strongest discrete sources re-overlaid (yellow ellipses)
- Black contours at 1.4, 1.9, 2.4, 2.9, 3.4, 3.9, 4.4, 10, 40K
- $\sigma \approx 250\text{mK}$ at 430 MHz

Region A (2 – 3 Mpc in extent) requires a distributed “fresh” energy source – plausibly provided by the ~ 7 embedded, radio galaxies.

Summary of signal detection

- In the hole zones, we have reached the absolute discrete source confusion limit.
- “diffuse” emission $\gtrsim 400$ mK is detected at many locations over the 70 sq deg field

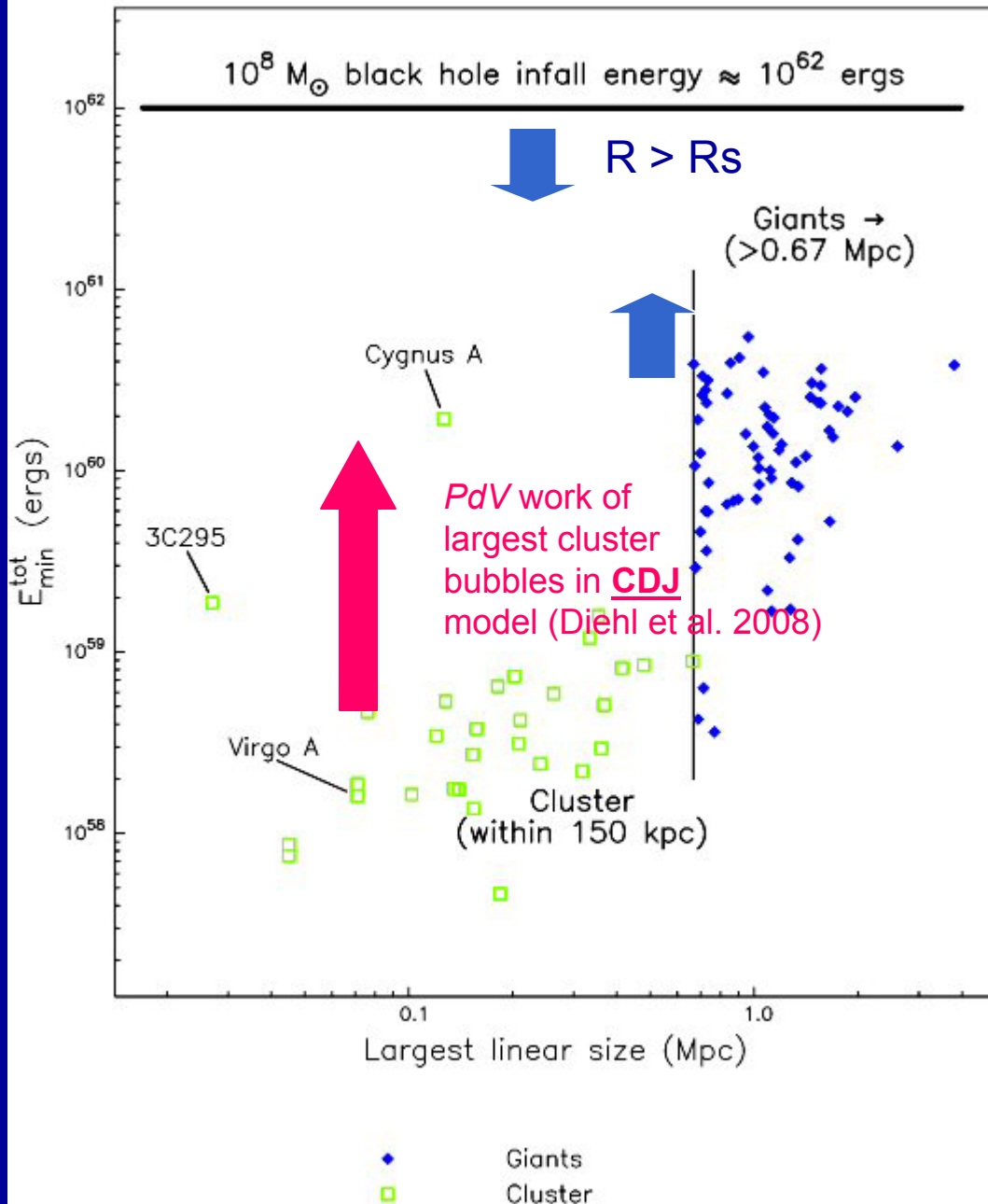
This is composed of:

1. Diffuse intergalactic emission (e.g. Regions A and B)
2. Galactic foreground (+ other extragalactic?), previously undetected on arcmin scales
3. Blends of faint discrete sources

5.

Energetics of intergalactic fields deriving from central BH's

Giant radio galaxies are the best calibrators of
BH energy input to the IGM (magnetic + CR)



$$= M_{\text{BH}} c^2$$

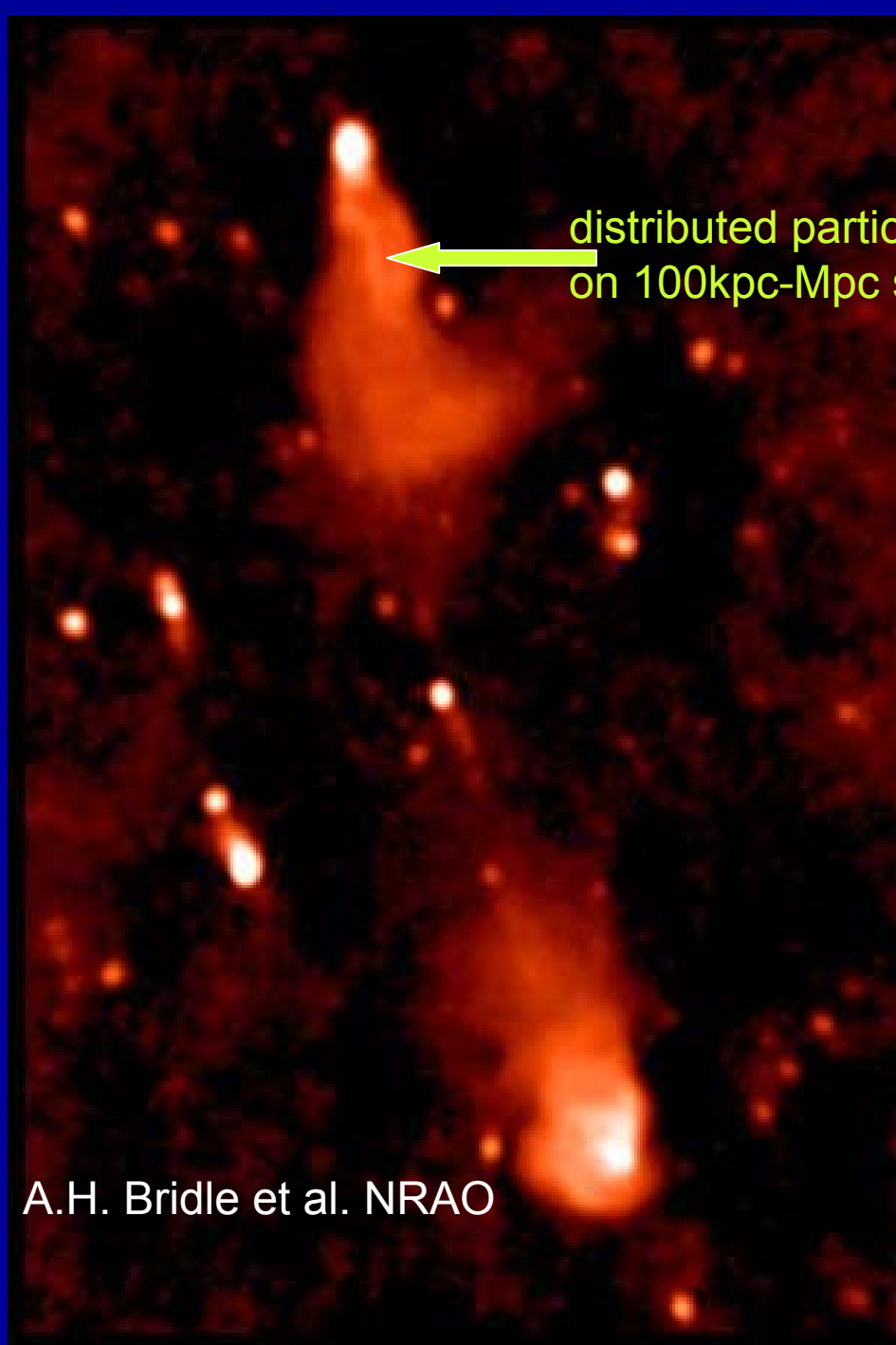


Mind the gap!!

Accumulated energy
($B^2/8\pi + \epsilon_{\text{CR}}$) \times (volume)
from "mature" BH-powered
radio source lobes

GRG's
capture the highest fraction
of the magnetic energy
released to the IGM

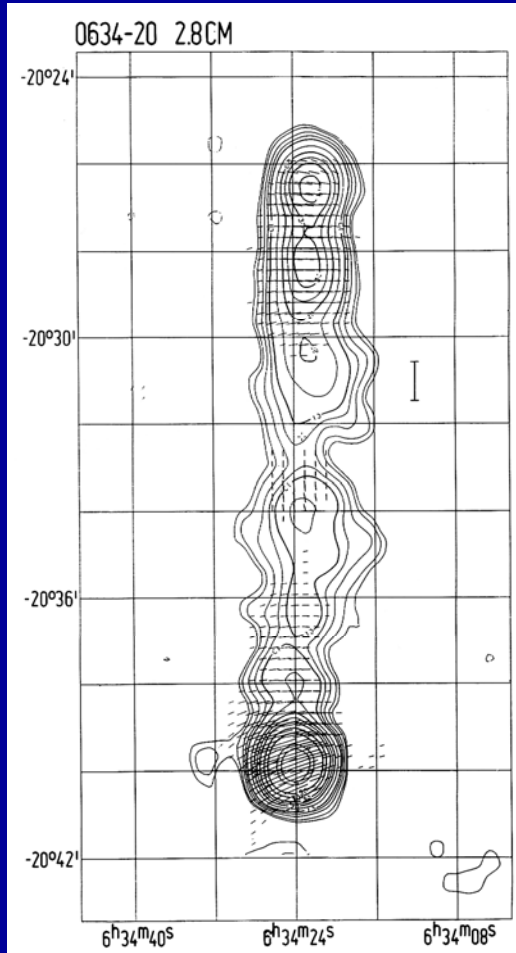
*Kronberg, Dufton, Li, &
Colgate,
ApJ 560, 178, 2001*



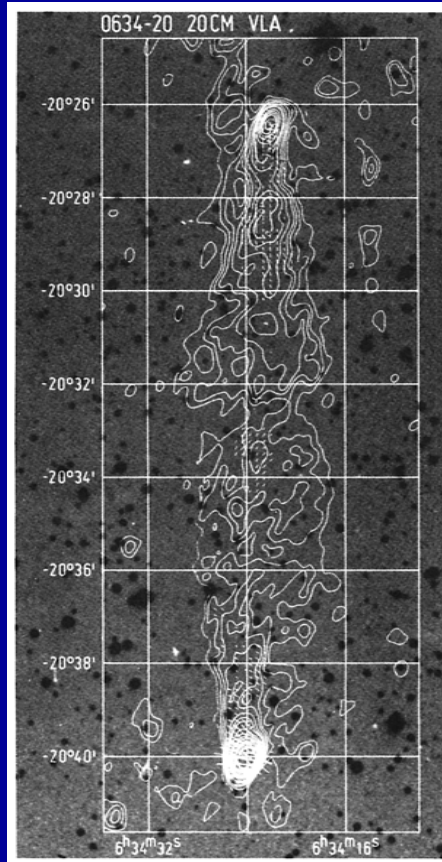
distributed particle acceleration
on 100kpc-Mpc scales

A.H. Bridle et al. NRAO

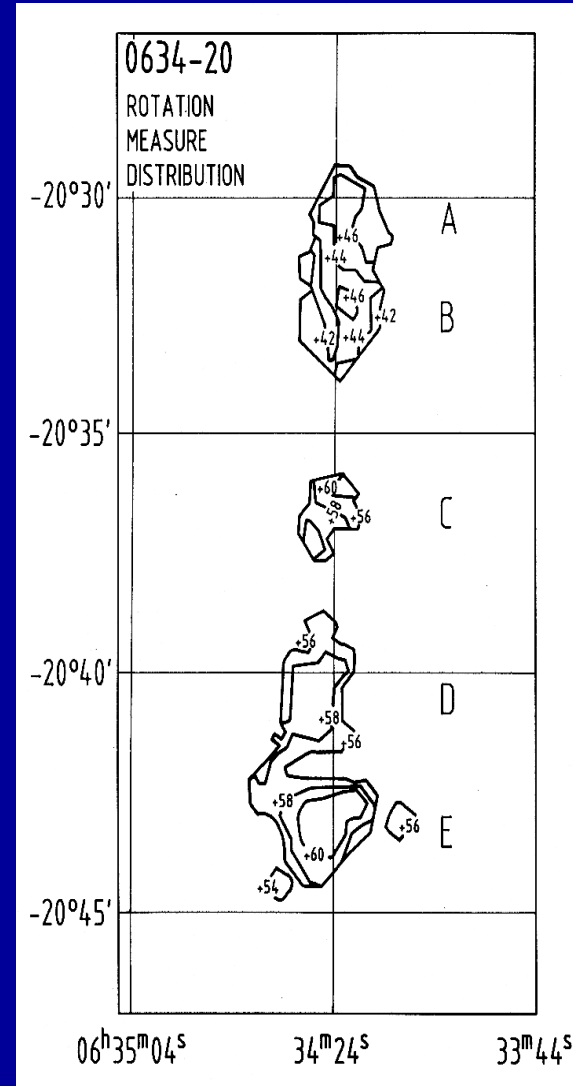
10 GHz



1.4 GHz



Faraday RM(radians/m²)



*Kronberg, Wielebinski & Graham
A&A 169, 63, 1986*

General properties of BH – fed radio lobes

- Thermal gas density is $<$ that of ambient IGM
- Easy to demonstrate within clusters,
- Appears also true of large radio lobes outside of clusters.
- Nicely demonstrated in WSRT GRG images of Strom & Willis 3C326, 3C236 A&A 1978, 1980, and others since

Opportunities for extending GRG observations with LOFAR and GMRT, incl. Faraday RM synthesis, at 1GHz \gtrsim f \gtrsim 150 MHz

6.

Evidence for dominance of magnetic structures in lobes and jets

- simulation of jet-lobe transition points in clusters
- jet/lobe systems as UHECR acceleration sites?
- cluster environment provide a ``controlled'' probe of lobe physics

Magnetic tower jet/lobe in a cluster environment

M. Nakamura, I.A. Tregillis, H. Li, S. Li ApJ 686 843, 2008

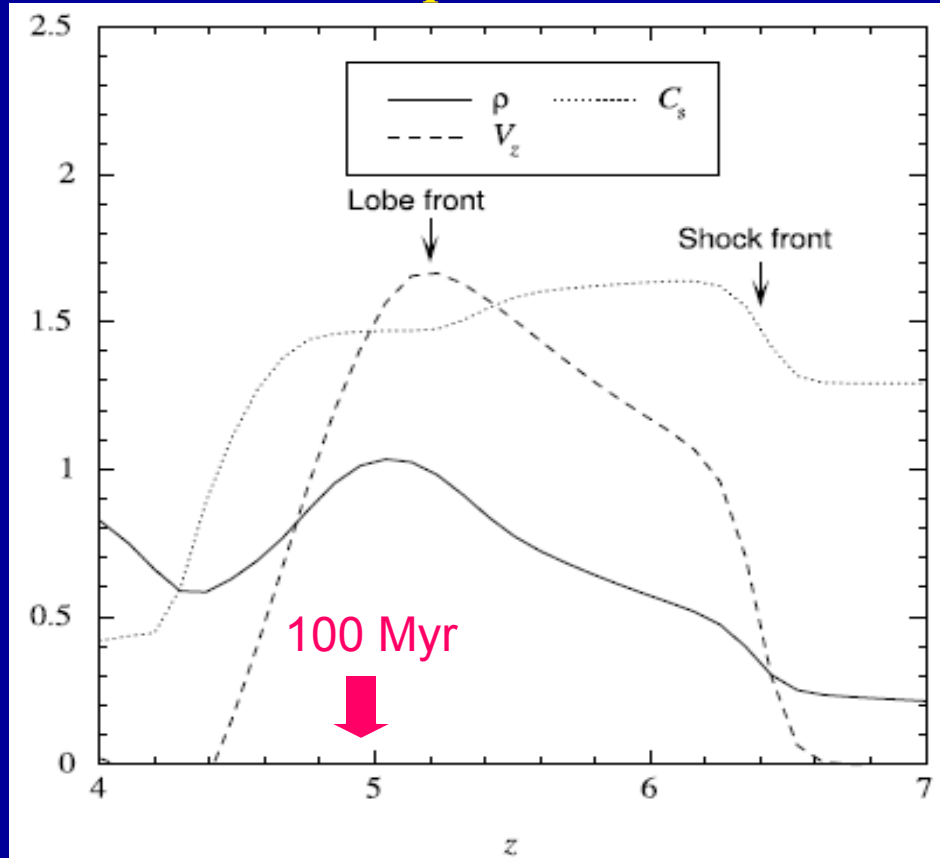
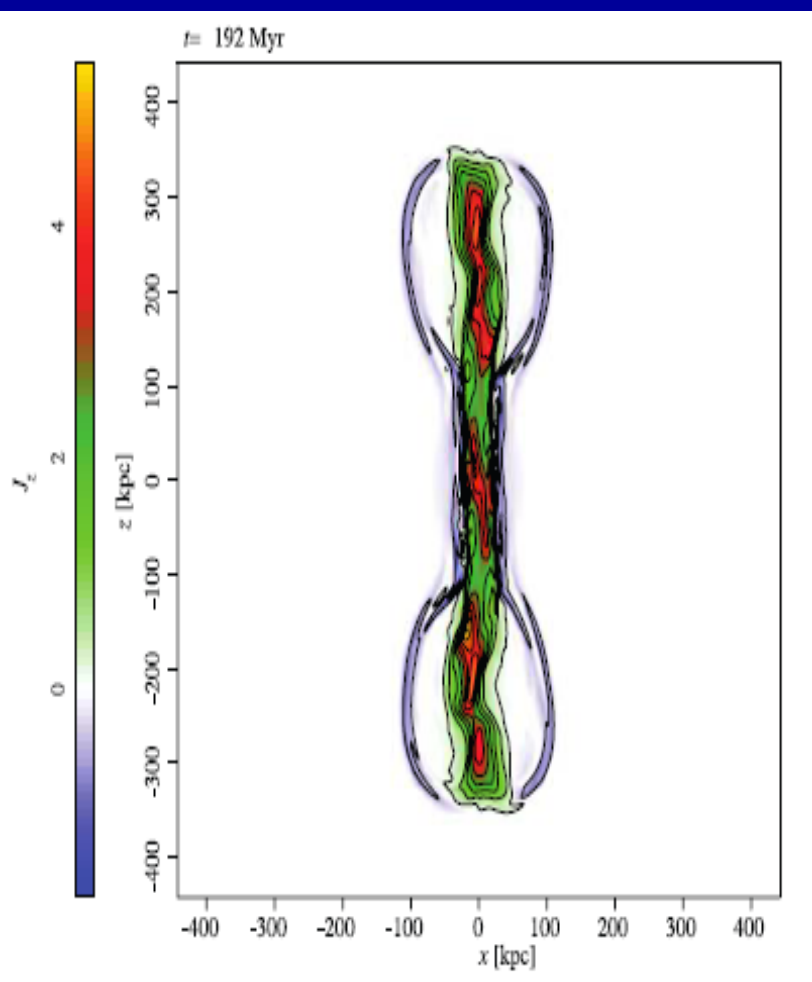
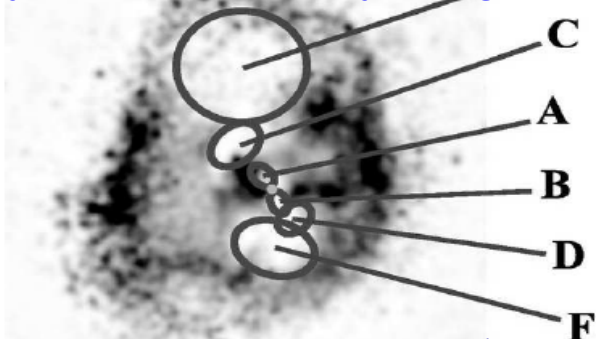


FIG. 2.— Axial profiles of physical quantities along the z -axis at $t = 3.0$ ($t = 72$ Myr): density ρ , sound speed C_s , and the axial velocity component V_z . The positions of the expanding shock and lobe fronts are shown.

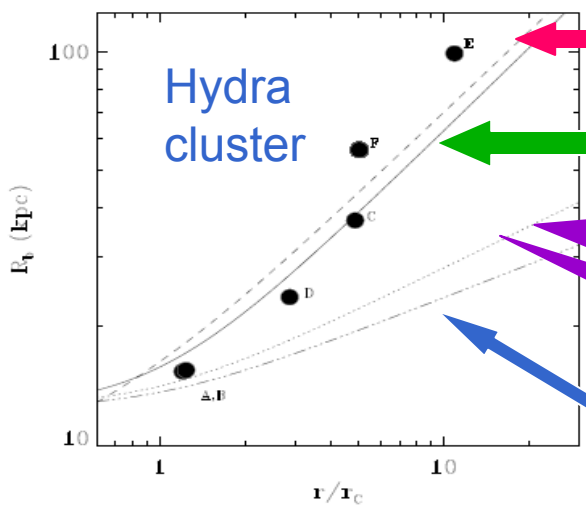
A recent galaxy cluster-environment test for
the
for the relative dominance of magnetic
BH/jet – energized lobes

Diehl, S., Li, H, Fryer, C, Rafferty, D ApJ 2008

Hydra cluster X-ray image:



Wise, M.W., McNamara, B.R., Nulsen, P.E.J, Houck, J.C., & David, L.P. ApJ 659, 1153, 2007



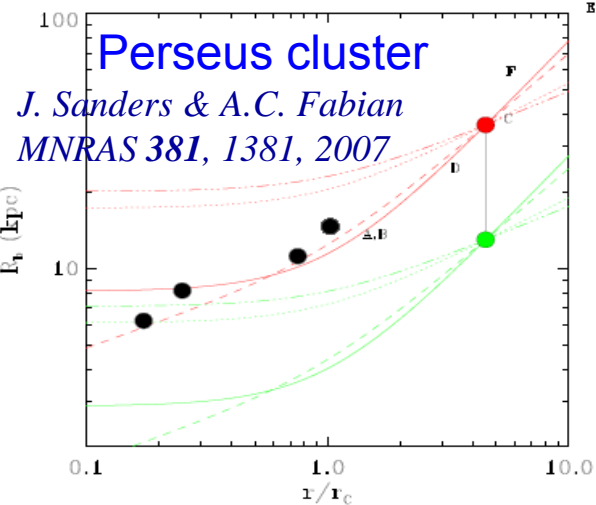
CIH continuous injection hydrodynamic model - - -

CDJ current-dominated MHD jet model ———

FML Bubbles contain frozen-in mag. loops

AD $\Gamma=5/3$, AD $\Gamma=4/3$ Adiabatically expanding hydrodynamic models

FIG. 6.— Left: The multi-cavity system in Hydra A, reproduced from Wise et al. (2007) with permission from the authors. The black area is excess X-ray emission left-over after an elliptical surface brightness model has been subtracted. Right: Data Points: Bubble sizes for Hydra A as a function of distance to the center, taken from Wise et al. (2007); Lines show predictions from the AD53 (triple-dot dashed line), AD43 (dotted line), FML (also dotted line), CIH (dashed line), as well as the CDJ model (solid line). The cavity labels are the same in both plots.



Perseus cluster
J. Sanders & A.C. Fabian
MNRAS 381, 1381, 2007

FIG. 7.— Bubble sizes for Perseus as a function of distance to the center. Lines as in Figure 6. The red data point shows the upper limit for the new bubble size estimate, the green data shows a lower limit. The correct answer will likely lie somewhere in between these two extremes.

limits to the true location of the bubbles. This will not only affect the radii themselves, but also the point at which other quantities are evaluated at, like density, temperature and pressure. In general the temperature rises outward in these systems, thus the temperature at the location of the bubble is likely to be systematically underestimated. The density and ambient pressure on the other hand will always be overestimated. This also means that any rise times derived from using the projected radius rather than the true distance to the center will result in estimates for the rise times that are systematically too low. We also note that the smaller the observed radius is, the higher the probability that it is due to an effect caused by projection.

But there are more subtle effects that projection has on our data. As we do not have an automated tool to detect bubbles, one has to rely on human experience in finding and identifying these systems. This task is much more difficult, if the cavities overlap with the bright cluster center or the bubble on the opposite side of the cluster. In fact, our sample does not contain any cavity system in which the bubble size exceeds the projected distance to the center, the slope of which is shown by the black solid line in Figure 8, even though this is statistically very improbable. This suggests that our sample is affected by what we will refer to as a “geometric” selection effect, introduced by our manual detection process.

Effects of Sig/Noise and projection effects;

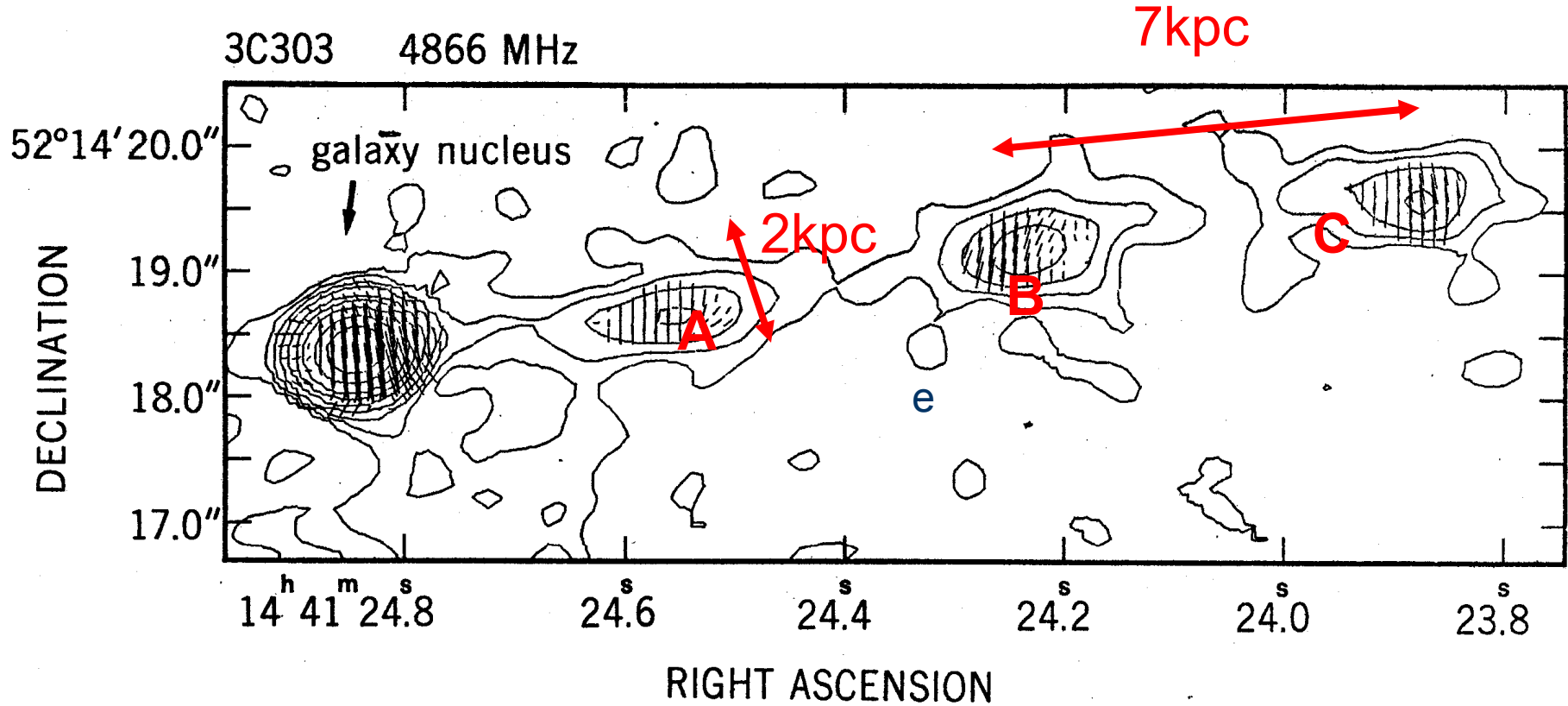
Enßlin & Heinz
A&A 384, L27, 2002

7.

kpc jets as potential sites for UHECR acceleration

- The 3C303 jet

Jets as UHECR accelerators?



$$E = \frac{B}{3 \text{ mG}} \times \frac{L}{1 \text{ kpc}} \Rightarrow 10^{19} \text{ eV}$$

Plasma Diagnostics of the 3C303 jet

Lapenta & Kronberg ApJ 625, 37-50, 2005

(1) $\langle \text{Energy flow rate} \rangle = E_{\text{min}}^T / \tau = 2.8 \times 10^{43} \tau_7^{-1} \text{ erg/s}$

(2) Total radio \rightarrow X-ray luminosity of the jet $= 1.7 \times 10^{42} \text{ erg s}^{-1}$

$$\frac{(2)}{(1)}$$

\rightarrow Radiative dissipation from the jet $\approx 10\%$ of energy flow rate along jet!

(3) Measure knots' synchrotron luminosity & size (D_{knot}) $\rightarrow B_{\text{int}}^{\text{knot}} = 10^{-3} \text{ G}$

(4) From the Faraday rotation images of the knots ($\text{RM} \propto n_{\text{th}} \times B_{\text{int}}^{\text{knot}} \times D_{\text{knot}}$)

$\rightarrow n_{\text{th}}$ in knots (upper limit for 3C303) $\rightarrow n_{\text{th}} \leq 1.4 \times 10^{-5} \text{ cm}^{-3}$

(3) & (4) \rightarrow lower limit to V_A within knots : $V_A^{\text{knot}} \propto B_{\text{int}}^{\text{knot}} / (n_{\text{th}})^{1/2}$

RESULT: $V_A^{\text{knot}} \approx 1.9c$ i.e. close to c

- For knot “C”, the RM image of 3C303 enables a measurement of the transverse ∇RM (radians/m²/m) over a knot. i.e. ∇RM is perpendicular to jet!
- B (RM) reverses sign on the jet axis. $|B|$ is estimated from measured synchrotron emissivity ($\gtrsim 1\text{mG}$)
- **a galaxy-scale, current-carrying “wire”**
- result for 3C303: $I = 7.5 \times 10^{17} (B_{-3}^G)$ [$r = 0.5\text{kpc}$]
ampères
- I is directed AWAY from the galaxy AGN nucleus in this knot
- Intrinsic knot polarization consistent with low- ϕ jet helical field

H. Ji, P.P. Kronberg, S.C. Prager, D. Uzdensky,
Physics of Plasmas **15**, 058302-8, 2008

8.

Background RM probes of magnetic fields in galaxy systems to high z ($\lesssim 6$)

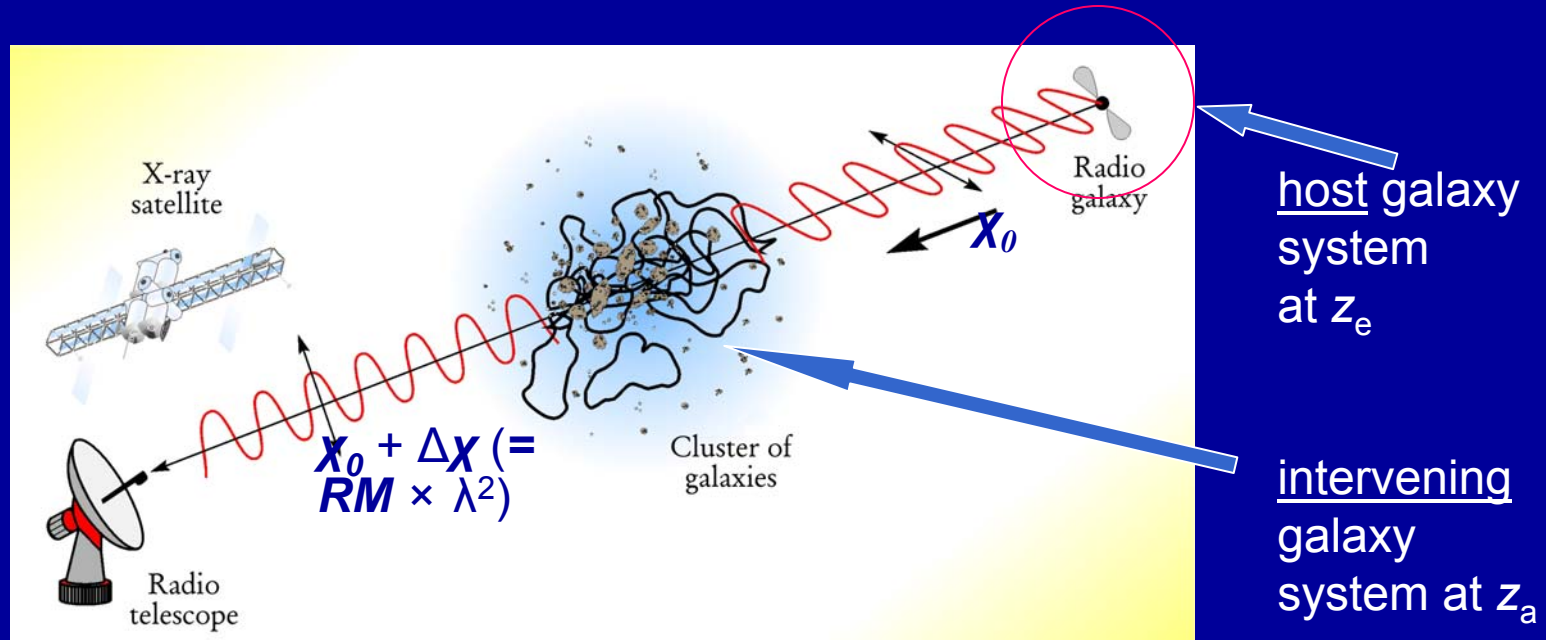
- Striking correlations exist between RM & high column density absorption systems in quasars, -- to large z .

Spectral resolution must be high enough to estimate W_{eq} . *Need 8+ meter optical telescopes to explore spectra to large z ! These now exist!*

Strong magnetic fields have been detected out to $z \gtrsim 3.5$! Slow (10^9 yr) galactic dynamos are not the explanation of $|B|$ amplification in galaxy systems.

- RM due to a widespread, co-expanding B_{igm} ?
No detections yet with current instruments

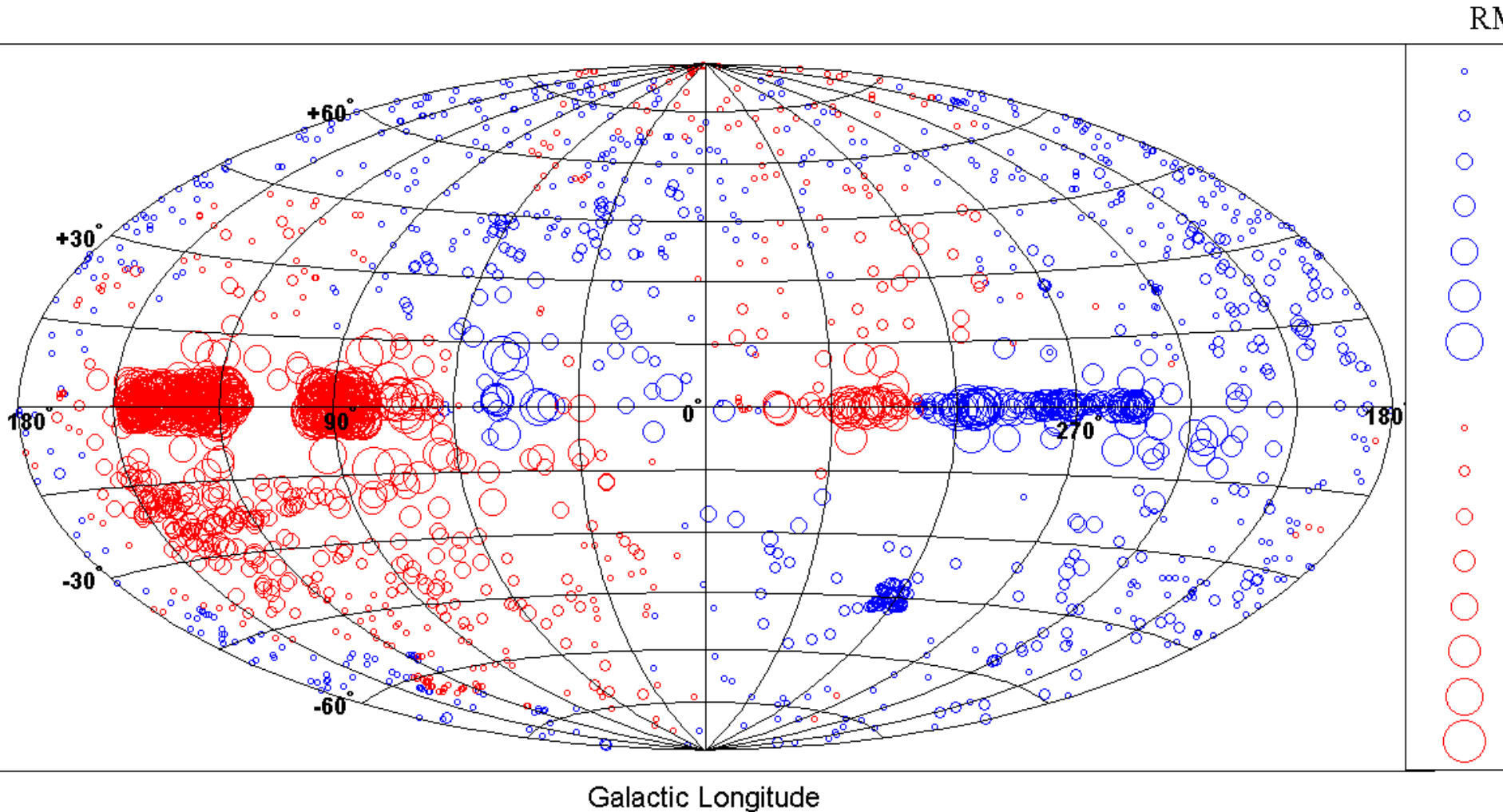
Faraday rotation at a distant EGRS, and at an intervenor



$$RM = \frac{\Delta\chi}{\Delta\lambda^2} = 8.12 \times 10^5 \int_0^{z_s} (1+z)^{-2} n_e(z) B_{\parallel}(z) dl(z) \quad \frac{\text{rad}}{\text{m}^2}$$

B in Gauss, n_e in cm^{-3} , l in pc

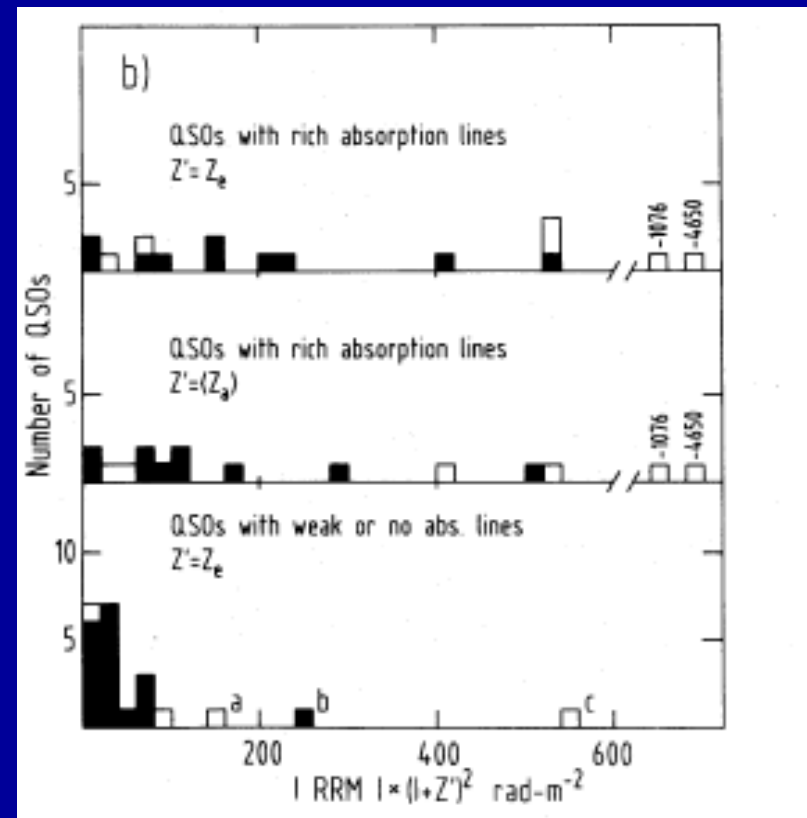
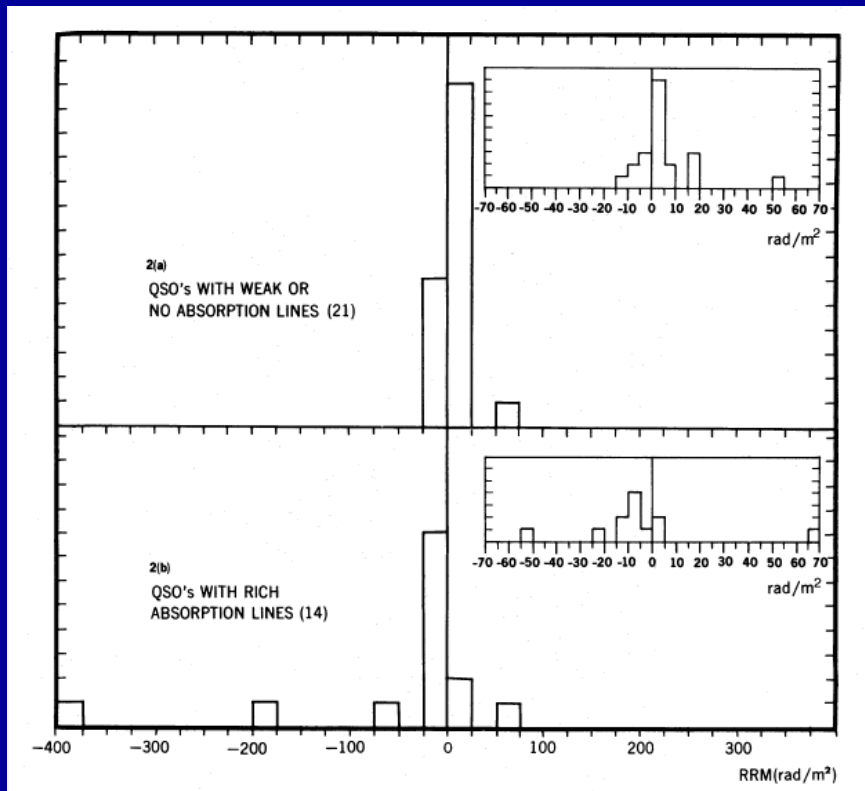
New smoothed Galactic RM sky from 2250 egrs RM's



Detections of magnetized optical absorption line systems

P.P. Kronberg & J.J. Perry,
ApJ 263, 518, 1982
(37 RM + Abs. spectrum QSO's)

G.L. Welter, J.J. Perry, & P.P. Kronberg
ApJ 279, 19, 1984
119 RM sample, 40 had spectra with
strong optical absorption lines

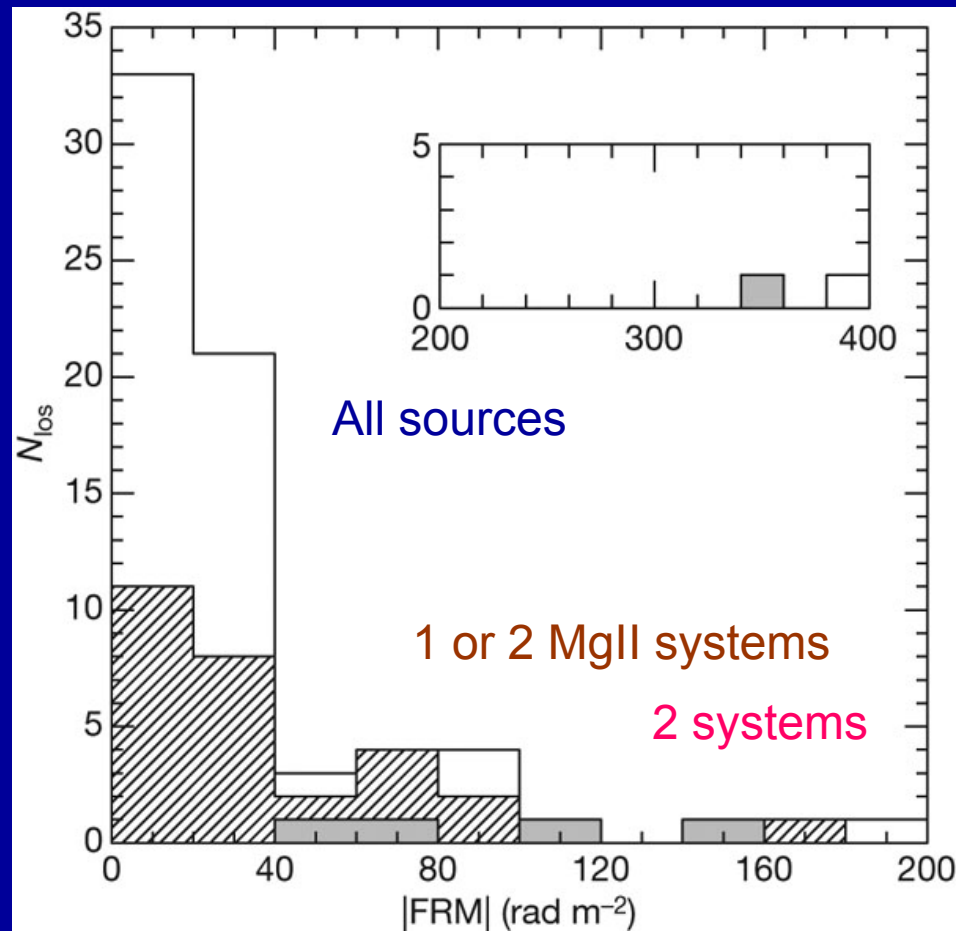


Effect of MgII absorption ($\lambda\lambda 2796.35, 2803.53 \text{ \AA}$) on the RM's of quasars

$2.0 \geq z \geq 0.6$ $m_V < 19$, and $|b| > 30^\circ$

M.L. Bernet, F. Miniati, S.J. Lilly, P.P. Kronberg, M. Dessauges-Zavadsky

Nature 454, 302-4, 2008 Jul18

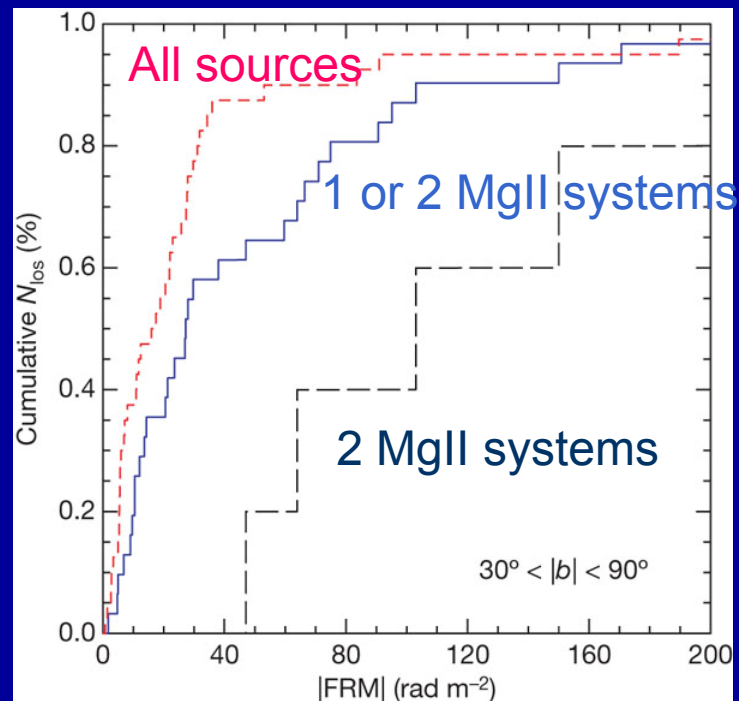


From new VLT
observations
in 2006-7.
UVES spectrograph

Cumulative plots of RM for 3 different MgII absorption line groups

M.L. Bernet, F. Miniati, S.J. Lilly, P.P. Kronberg, M. Dessauges-Zavadsky Nature 454, 302-4, 2008

Method: G.L. Welter, J.J. Perry & P.P. Kronberg ApJ 279, 19, 1984

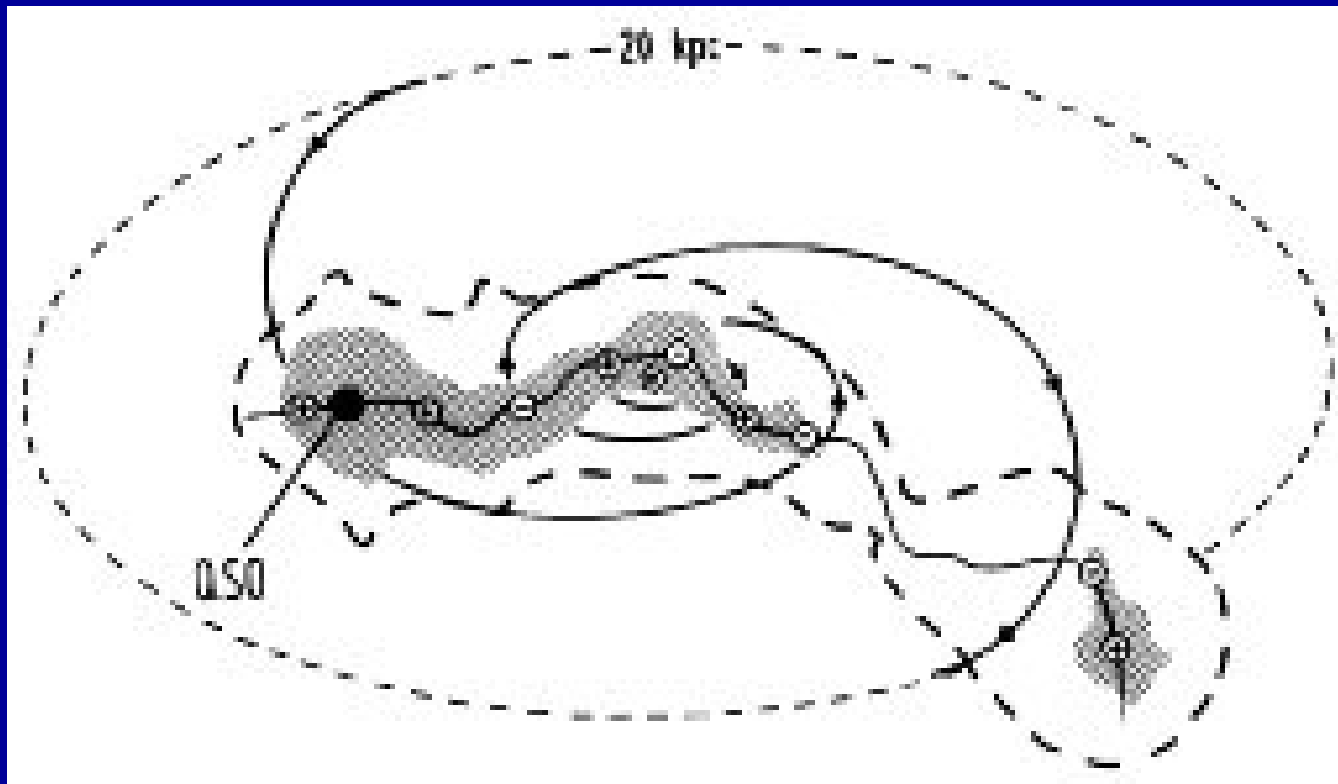


2-D magnetic probes of intervenors at high z

- High- z jets make promising probes of intervening galaxy systems.
- 2 systems described below
 - (1) – an intervening spiral-like galaxy ($z = 0.38$) in front of a $z = 1$ quasar jet
 - (2) -- an “associated” absorbing gas cloud at $z_{\text{abs}} \approx z_{\text{emission}}$ for 3C191 ($z \sim 1.9$)
- **Future: 3-D** using RM synthesis techniques

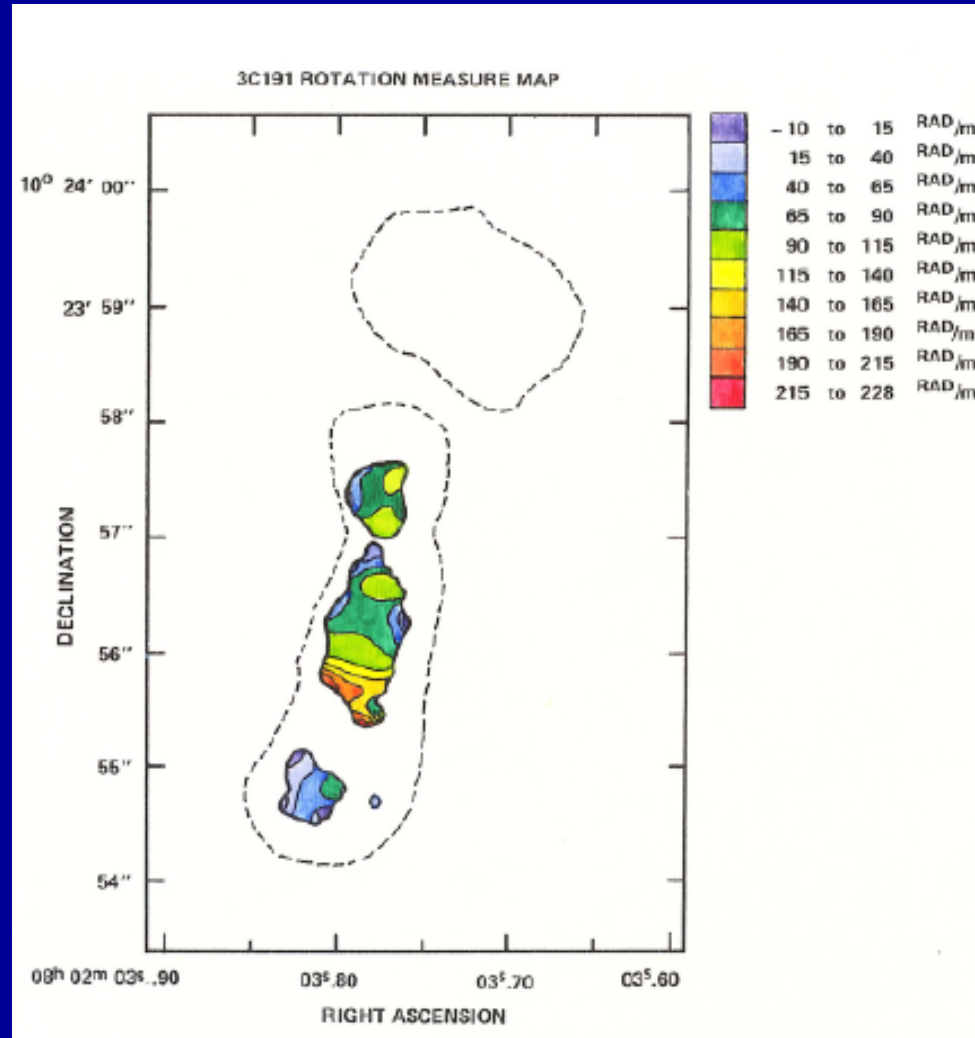
PKS 1229-021 a jet quasar at $z = 1.03$ behind a
spiral galaxy at $z = 0.4$

Kronberg, Perry & Zukowski ApJ 387, 528-535, 1992



3C191 a quasar jet ($z = 1.95$) with an “associated” intervenor with rich absorption lines

*P.P Kronberg,
J.J.Perry &
E.L.H. Zukowski
ApJL 355, L31, 1990*



RM
measured
between
5GHz and 15
GHz

Principle conclusion from both analyses of (1) all sources, and (2) MgII absorbers

- Magnetic field strengths in galaxy systems ($N_e \gtrsim 10^{20} \text{ cm}^{-2}$) up to $\gtrsim 80\%$ of a Hubble time ago are at least comparable to those at $z = 0$.
- *i.e.* confirms lack of evidence for a slow galactic dynamo field amplification over cosmic time

10.

Can we detect $|\mathbf{B}| \neq 0$ before the epoch of recombination?

Currently, few possibilities to detect \mathbf{B} at $\tau < \tau_{\text{RECOMB}}$.

Best(?) one: look for a Faraday RM signal in the polarized CMB over the appropriate range of multipole scales (l). (**Kosowsky & Loeb ApJ 469, 1 1996**)

$$RM = \Delta\chi / (\Delta\lambda)^2 \text{ (radians/m}^2\text{)} = k \int n_e B_{\parallel} dl$$

- $\Delta\chi$ must be a detectable angle rotation
- Longest λ in $\Delta\lambda$ range must be short enough to be free of all polarized foregrounds at $z \lesssim 1000$
- Need to evolve n_e profile as scattering τ increases through the recombination redshift.

(W. Hu, D. Scott, N. Sugiyama & M. White, Phys Rev. D 52, 5498, 1995)

Illustrative result

$$\langle \Delta\chi_{\nu_1, \nu_2}^2 \rangle^{1/2} = 1.1^\circ \left(1 - \frac{\nu_1^2}{\nu_2^2} \right) \left(\frac{B_0}{10^{-9} G} \right) \left(\frac{30 \text{ GHz}}{\nu_1} \right)$$

Kosowsky & Loeb ApJ 469, 1 1996

i.e. 1.1° rotation at 30 GHz (λ 1cm) \Rightarrow RM 195 rad m^{-2} at $z = 0$

ν must be *high enough to be beyond* all foreground polarized rad'n
(e.g. synchrotron, polarized dust, etc. at any $z \lesssim 1000$)

ν must be *low enough* to detect rotation

$\nu_1 \sim 30$ GHz is \sim the only possible window, and even there $\Delta\chi$ is only $\sim 1^\circ$

A direct, but difficult measurement. May be possible in future

11.

How to detect Magnetic fields in cosmic voids?

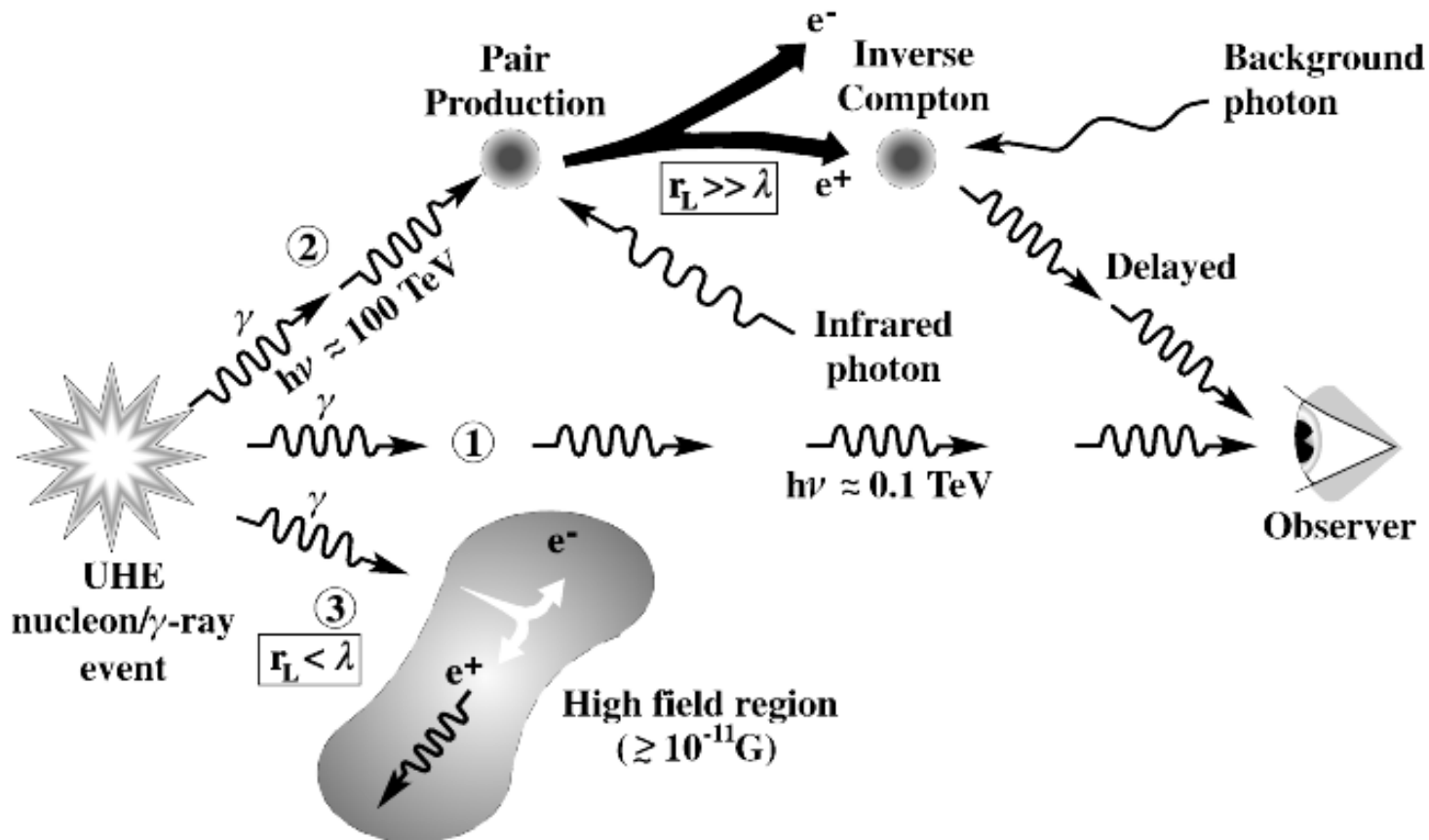
- Diffusion out of the walls and filaments?
- Relic of a pre-galactic field?

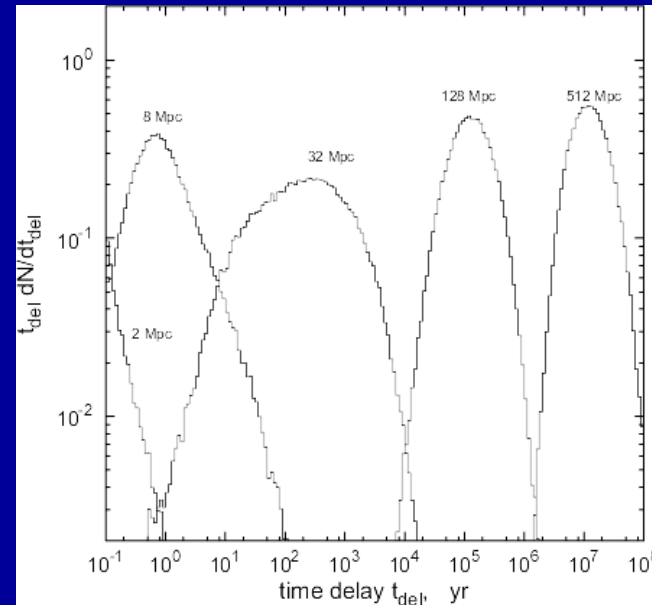
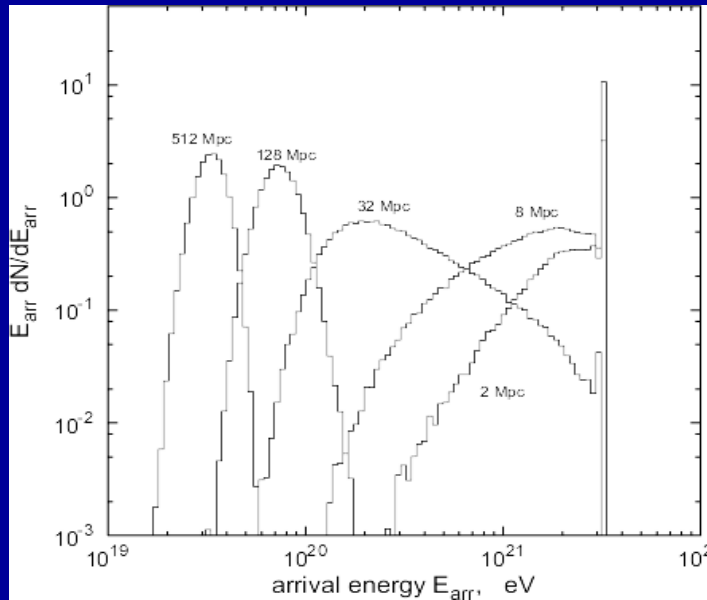
(these two need to be independently verified)

- Propagation of both high energy particles & photons
- Time of arrival, deflection, energy and composition

Energy dependent cascade of a broadband γ -ray burst might probe a very weak IGM field

High energy $h\nu - e^+e^-$ cascades in the intergalactic medium





Stanev T., Engel, R., Mücke, A., Protheroe, R. J., Rachen, J. Phys Rev. D, 62, 0930052000

(left)

The received CR energy distribution on Earth for a monoenergetically injected proton energy of $10^{21.5}$ eV for a randomly orientated $B_{IG} = 10^{-9}$ G at progressively larger distances, up to 512 Mpc. The energy is reduced by the GZK effect (most severe), B-H pair production losses, and adiabatic losses.

(right)

The relative time delay for protons injected **at the same distance**, when propagated through a randomly oriented magnetic field of 10^{-9} G, where $l_0 = 1$ Mpc.

Limitations in current observational diagnostics = opportunities!

- Need better resolution and frequency coverage for AGN jet and lobe RM images – for synergy w. simulations
- Deeper X-ray observations of jets ($\lesssim 1''$ res.)
- X-ray and EUV observations in the IGM
- γ - ray observations
- Milky Way foregrounds in more deep synchrotron surveys
- More discrete source Faraday rotation observations needed

End

P.P. Kronberg